

Fully discrete polynomial de Rham sequences of arbitrary degree on polyhedral meshes

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References

- **Local DDR sequences** [DP, Droniou and Rapetti, 2020]
- **Global DDR sequences** and stability [DP and Droniou, 2020a]
- Primal and dual **consistency** [DP and Droniou, ongoing]
- See [DP and Droniou, 2020b] for polytopal analysis tools



Outline

- 1 Introduction and motivation
- 2 Discrete de Rham (DDR) sequences
- 3 Application to magnetostatics

A (not so simple) model problem I

- Let $\Omega \subset \mathbb{R}^3$ be an open connected polyhedral domain that **does not enclose any void**
- Let a **current density** $\mathbf{f} \in \mathbf{curl} \mathbf{H}(\mathbf{curl}; \Omega)$ be given
- We consider the problem: Find the **magnetic field** $\boldsymbol{\sigma} : \Omega \rightarrow \mathbb{R}^3$ and the **vector potential** $\mathbf{u} : \Omega \rightarrow \mathbb{R}^3$ s.t.

$$\boldsymbol{\sigma} - \mathbf{curl} \mathbf{u} = \mathbf{0} \quad \text{in } \Omega, \quad (\text{vector potential})$$

$$\mathbf{curl} \boldsymbol{\sigma} = \mathbf{f} \quad \text{in } \Omega, \quad (\text{Ampère's law})$$

$$\operatorname{div} \mathbf{u} = 0 \quad \text{in } \Omega, \quad (\text{Coulomb's gauge})$$

$$\mathbf{u} \times \mathbf{n} = \mathbf{0} \quad \text{on } \partial\Omega \quad (\text{boundary condition})$$

- The extension to variable magnetic permeability is straightforward

A (not so simple) model problem II

- In **weak formulation**: Find $(\boldsymbol{\sigma}, \mathbf{u}) \in \mathbf{H}(\mathbf{curl}; \Omega) \times \mathbf{H}(\mathbf{div}; \Omega)$ s.t.

$$\begin{aligned} \int_{\Omega} \boldsymbol{\sigma} \cdot \boldsymbol{\tau} - \int_{\Omega} \mathbf{u} \cdot \mathbf{curl} \boldsymbol{\tau} &= 0 & \forall \boldsymbol{\tau} \in \mathbf{H}(\mathbf{curl}; \Omega), \\ \int_{\Omega} \mathbf{curl} \boldsymbol{\sigma} \cdot \boldsymbol{\nu} + \int_{\Omega} \mathbf{div} \mathbf{u} \mathbf{div} \boldsymbol{\nu} &= \int_{\Omega} \mathbf{f} \cdot \boldsymbol{\nu} & \forall \boldsymbol{\nu} \in \mathbf{H}(\mathbf{div}; \Omega) \end{aligned}$$

- **Well-posedness** hinges on the **exactness** of the following portion of the de Rham sequence:

$$\mathbb{R} \xrightarrow{i_{\Omega}} H^1(\Omega) \xrightarrow{\text{grad}} \mathbf{H}(\mathbf{curl}; \Omega) \xrightarrow{\mathbf{curl}} \mathbf{H}(\mathbf{div}; \Omega) \xrightarrow{\mathbf{div}} L^2(\Omega) \xrightarrow{0} \{0\}$$

- **This exactness property is also needed at the discrete level!**

The Finite Element way

Local spaces

- **Key idea:** define **subspaces** that form **exact sequence**
- Let $T \subset \mathbb{R}^3$ be a polyhedron and set, for any $k \geq -1$,

$$\mathcal{P}^k(T) := \{\text{restrictions of 3-variate polynomials of degree } \leq k \text{ to } T\}$$

- Fix $k \geq 0$ and write, denoting by \mathbf{x}_T the barycenter of T ,

$$\begin{aligned}\mathcal{P}^k(T)^3 &= \underbrace{\text{grad } \mathcal{P}^{k+1}(T)}_{\mathcal{G}^k(T)} \oplus \underbrace{(\mathbf{x} - \mathbf{x}_T) \times \mathcal{P}^{k-1}(T)^3}_{\mathcal{G}^{c,k}(T)} \\ &= \underbrace{\text{curl } \mathcal{P}^{k+1}(T)^3}_{\mathcal{R}^k(T)} \oplus \underbrace{(\mathbf{x} - \mathbf{x}_T) \mathcal{P}^{k-1}(T)}_{\mathcal{R}^{c,k}(T)}\end{aligned}$$

- Define the **trimmed spaces**

$$\mathcal{N}^k(T) := \mathcal{G}^k(T) \oplus \mathcal{G}^{c,k+1}(T) \quad [\text{Nédélec, 1980}]$$

$$\mathcal{RT}^k(T) := \mathcal{R}^k(T) \oplus \mathcal{R}^{c,k+1}(T) \quad [\text{Raviart and Thomas, 1977}]$$

The Finite Element way

Global FE sequence

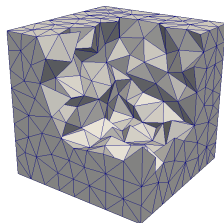


Figure: Conforming tetrahedral mesh of the unit cube (clip)

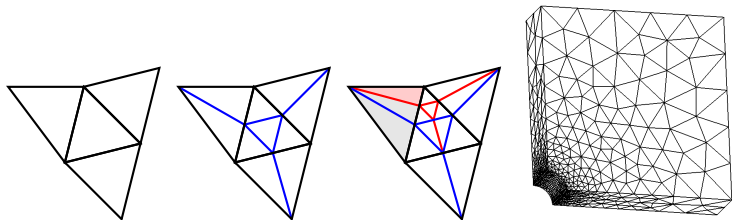
- Let $\mathcal{T}_h = \{T\}$ be a **conforming tetrahedral mesh** of Ω and let $k \geq 0$
- Local spaces can be **glued together** to form the **global FE sequence**

$$\mathbb{R} \xrightarrow{i_\Omega} \mathcal{P}_c^{k+1}(\mathcal{T}_h) \xrightarrow{\text{grad}} \mathcal{N}^k(\mathcal{T}_h) \xrightarrow{\text{curl}} \mathcal{RT}^k(\mathcal{T}_h) \xrightarrow{\text{div}} \mathcal{P}^k(\mathcal{T}_h) \xrightarrow{0} \{0\}$$

- **This procedure only works on conforming meshes!**

The Finite Element way

Shortcomings



- Approach limited to **conforming meshes** with **standard elements**
- \implies local refinement requires to **trade mesh size for mesh quality**
- \implies complex geometries may require a **large number of elements**
- \implies the element shape cannot be **adapted to the solution**
- The implementation of **high-order** versions may be tricky
- ...

The discrete de Rham (DDR) approach I

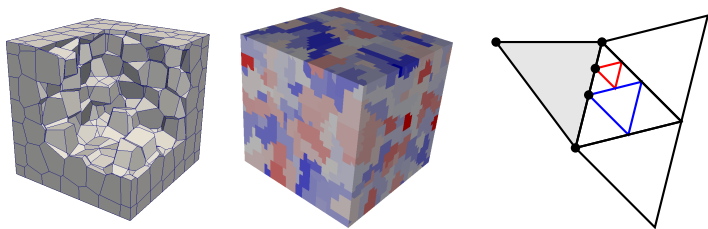


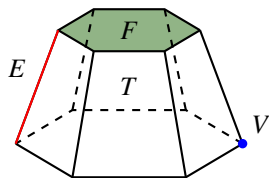
Figure: Examples of polytopal meshes supported by the DDR approach

- **Key idea:** replace spaces **and operators** by discrete counterparts:

$$\mathbb{R} \xrightarrow{I_{\text{grad},h}^k} \underline{X}_{\text{grad},h}^k \xrightarrow{\underline{G}_h^k} \underline{X}_{\text{curl},h}^k \xrightarrow{\underline{C}_h^k} \underline{X}_{\text{div},h}^k \xrightarrow{D_h^k} \mathcal{P}^k(\mathcal{T}_h) \xrightarrow{0} \{0\}$$

- Support of **general polyhedral meshes** and **high-order (!)**
- Exactness proved **at the discrete level** (directly usable for stability)
- (Relatively) simple implementation of **high-order versions**

The discrete de Rham (DDR) approach II



- The fully discrete spaces are spanned by **vectors of polynomials**
- Polynomial components **attached to geometric objects** to mimic
 - **full continuity** for the approximation of $H^1(\Omega)$
 - **continuity of tangential traces** for the approximation of $\mathbf{H}(\text{curl}; \Omega)$
 - **continuity of normal traces** for the approximation of $\mathbf{H}(\text{div}; \Omega)$
- Selected so as to enable the reconstruction of consistent
 - discrete **vector calculus operators**
 - (scalar or vector) **discrete potentials**

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The two-dimensional case

Continuous exact sequence

- Let $F \subset \mathbb{R}^2$ be a **simply connected polygon** embedded in \mathbb{R}^3
- Let, for $q : F \rightarrow \mathbb{R}$ and $\mathbf{v} : F \rightarrow \mathbb{R}^2$ smooth enough,

$$\mathbf{rot}_F q := \varrho_{-\pi/2}(\mathbf{grad}_F q) \quad \mathbf{rot}_F \mathbf{v} := \operatorname{div}_F(\varrho_{-\pi/2}\mathbf{v})$$

- We derive a discrete counterpart of the **exact local sequence**:

$$\mathbb{R} \xrightarrow{i_F} H^1(F) \xrightarrow{\mathbf{grad}_F} \mathbf{H}(\mathbf{rot}; F) \xrightarrow{\mathbf{rot}_F} L^2(F) \xrightarrow{0} \{0\}$$

- We will need the following decompositions of $\mathcal{P}^k(F)^2$:

$$\begin{aligned} \mathcal{P}^k(F)^2 &= \underbrace{\mathbf{rot}_F \mathcal{P}^{k+1}(F)}_{\mathcal{R}^k(F)} \oplus \underbrace{(\mathbf{x} - \mathbf{x}_F)\mathcal{P}^{k-1}(F)}_{\mathcal{R}^{c,k}(F)} \\ &= \underbrace{\mathbf{grad}_F \mathcal{P}^{k+1}(F)}_{\mathcal{G}^k(F)} \oplus \underbrace{(\mathbf{x} - \mathbf{x}_F)^\perp \mathcal{P}^{k-1}(F)}_{\mathcal{G}^{c,k}(F)} \end{aligned}$$

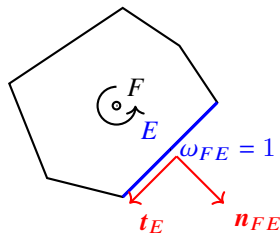
The two-dimensional case

A key remark

- Denote by $\pi_{\mathcal{P},F}^{k-1}$ the L^2 -orthogonal projector on $\mathcal{P}^{k-1}(F)$
- Let $q \in \mathcal{P}^{k+1}(F)$. For any $\mathbf{v} \in \mathcal{P}^k(F)^2$, we have

$$\begin{aligned}\int_F \mathbf{grad}_F q \cdot \mathbf{v} &= - \int_F q \underbrace{\operatorname{div}_F \mathbf{v}}_{\in \mathcal{P}^{k-1}(F)} + \sum_{E \in \mathcal{E}_F} \omega_{FE} \int_E q|_{\partial F} (\mathbf{v} \cdot \mathbf{n}_{FE}) \\ &= - \int_F \pi_{\mathcal{P},F}^{k-1} q \operatorname{div}_F \mathbf{v} + \sum_{E \in \mathcal{E}_F} \omega_{FE} \int_E q|_{\partial F} (\mathbf{v} \cdot \mathbf{n}_{FE})\end{aligned}$$

- Hence, $\mathbf{grad}_F q$ can be computed given $\pi_{\mathcal{P},F}^{k-1} q$ and $q|_{\partial F}$



The two-dimensional case

Discrete $H^1(F)$ space

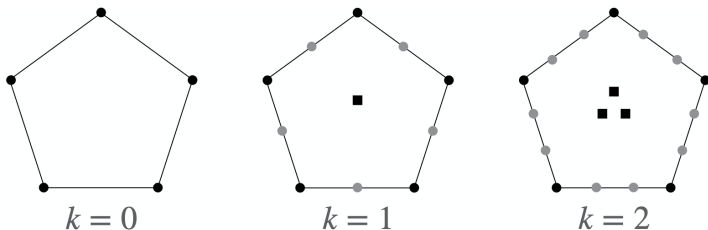


Figure: Number of degrees of freedom for $\underline{X}_{\text{grad},F}^k$ for $k \in \{0, 1, 2\}$

- The discrete counterpart of $H^1(F)$ is

$$\underline{X}_{\text{grad},F}^k := \left\{ \underline{q}_F = (q_F, q_{\partial F}) : q_F \in \mathcal{P}^{k-1}(F) \text{ and } q_{\partial F} \in \mathcal{P}_c^{k+1}(\mathcal{E}_F) \right\}$$

- The interpolator $\underline{I}_{\text{grad},F}^k : C^0(\bar{F}) \rightarrow \underline{X}_{\text{grad},F}^k$ is s.t., $\forall q \in C^0(\bar{F})$,

$$\underline{I}_{\text{grad},F}^k q := (\pi_{\mathcal{P},F}^{k-1} q, q_{\partial F}) \text{ with}$$

$$\pi_{\mathcal{P},E}^{k-1}(q_{\partial F})|_E = \pi_{\mathcal{P},E}^{k-1} q|_E \quad \forall E \in \mathcal{E}_F \text{ and } q_{\partial F}(\mathbf{x}_V) = q(\mathbf{x}_V) \quad \forall V \in \mathcal{V}_F$$

The two-dimensional case

Reconstructions in $\underline{X}_{\text{grad},F}^k$

- For all $E \in \mathcal{E}_F$, the **edge gradient** $G_E^k : \underline{X}_{\text{grad},F}^k \rightarrow \mathcal{P}^k(E)$ is s.t.

$$G_E^k q_{\underline{F}} := (q \partial F)'|_E$$

- The **full face gradient** $G_F^k : \underline{X}_{\text{grad},F}^k \rightarrow \mathcal{P}^k(F)^2$ is s.t., $\forall \mathbf{v} \in \mathcal{P}^k(F)^2$,

$$\int_F G_F^k q_{\underline{F}} \cdot \mathbf{v} = - \int_F q_F \operatorname{div}_F \mathbf{v} + \sum_{E \in \mathcal{E}_F} \omega_{FE} \int_E q \partial F (\mathbf{v} \cdot \mathbf{n}_{FE})$$

- By construction, we have **polynomial consistency**:

$$G_F^k (\underline{I}_{\text{grad},F}^k q) = \mathbf{grad}_F q \quad \forall q \in \mathcal{P}^{k+1}(F)$$

- We reconstruct similarly a **face potential** in $\mathcal{P}^{k+1}(F)$

The two-dimensional case

Discrete $\mathbf{H}(\text{rot}; F)$ space

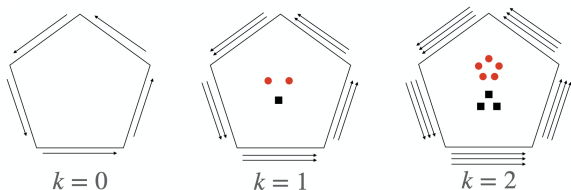


Figure: Number of degrees of freedom for $\underline{\mathbf{X}}_{\text{rot},F}^k$ for $k \in \{0, 1, 2\}$

- We reason starting from: $\forall \mathbf{v} \in \mathcal{N}^k(F) := \mathcal{G}^k(F) \oplus \mathcal{G}^{c,k+1}(F)$,

$$\int_F \text{rot}_F \mathbf{v} q = \int_F \mathbf{v} \cdot \underbrace{\text{rot}_F q}_{\in \mathcal{R}^{k-1}(F)} - \sum_{E \in \mathcal{E}_F} \omega_{FE} \int_E \underbrace{(\mathbf{v} \cdot \mathbf{t}_E)}_{\in \mathcal{P}^k(E)} q|_E \quad \forall q \in \mathcal{P}^k(F)$$

- This leads to the following discrete counterpart of $\mathbf{H}(\text{rot}; F)$:

$$\underline{\mathbf{X}}_{\text{rot},F}^k := \left\{ \mathbf{v}_F = (\mathbf{v}_{\mathcal{R},F}, \mathbf{v}_{\mathcal{R},F}^c, (\mathbf{v}_E)_{E \in \mathcal{E}_F}) : \right.$$

$$\left. \mathbf{v}_{\mathcal{R},F} \in \mathcal{R}^{k-1}(F), \mathbf{v}_{\mathcal{R},F}^c \in \mathcal{R}^{c,k}(F), \mathbf{v}_E \in \mathcal{P}^k(E) \forall E \in \mathcal{E}_F \right\}$$

The two-dimensional case

Reconstructions in $\underline{\mathbf{X}}_{\text{rot},F}^k$

- The **face curl operator** $C_F^k : \underline{\mathbf{X}}_{\text{rot},F}^k \rightarrow \mathcal{P}^k(F)$ is s.t.,

$$\int_F C_F^k \underline{\mathbf{v}}_F q = \int_F \mathbf{v}_{\mathcal{R},F} \cdot \mathbf{rot}_F q - \sum_{E \in \mathcal{E}_F} \omega_{FE} \int_E \mathbf{v}_E q \quad \forall q \in \mathcal{P}^k(F)$$

- Define the interpolator $\underline{\mathbf{I}}_{\text{rot},F}^k : H^1(F)^2 \rightarrow \underline{\mathbf{X}}_{\text{rot},F}^k$ s.t., $\forall \mathbf{v} \in H^1(F)^2$,

$$\underline{\mathbf{I}}_{\text{rot},F}^k \mathbf{v} := \left(\boldsymbol{\pi}_{\mathcal{R},F}^{k-1} \mathbf{v}, \boldsymbol{\pi}_{\mathcal{R},F}^{c,k} \mathbf{v}, \left(\boldsymbol{\pi}_{\mathcal{P},E}^k (\mathbf{v}|_E \cdot \mathbf{t}_E) \right)_{E \in \mathcal{E}_F} \right).$$

- C_F^k is **polynomially consistent** by construction:

$$C_F^k (\underline{\mathbf{I}}_{\text{rot},F}^k \mathbf{v}) = \mathbf{rot}_F \mathbf{v} \quad \forall \mathbf{v} \in \mathcal{N}^k(F)$$

- By similar principles, we reconstruct a **vector potential** in $\mathcal{P}^k(F)^2$

The two-dimensional case

Exact local sequence

Theorem (Exactness of the two-dimensional local DDR sequence)

If F is simply connected, the following local sequence is *exact*:

$$\mathbb{R} \xrightarrow{I_{\text{grad},F}^k} \underline{X}_{\text{grad},F}^k \xrightarrow{\underline{G}_F^k} \underline{X}_{\text{rot},F}^k \xrightarrow{C_F^k} \mathcal{P}^k(F) \xrightarrow{0} \{0\},$$

where $\underline{G}_F^k : \underline{X}_{\text{grad},F}^k \rightarrow \underline{X}_{\text{rot},F}^k$ is the *discrete gradient* s.t., $\forall \underline{q}_F \in \underline{X}_{\text{grad},F}^k$,

$$\underline{G}_F^k \underline{q}_F := \left(\pi_{\mathcal{R},F}^{k-1} (G_F^k \underline{q}_F), \pi_{\mathcal{R},F}^{c,k} (G_F^k \underline{q}_F), (G_E^k \underline{q}_F)_{E \in \mathcal{E}_F} \right)$$

The two-dimensional case

Summary

$$\mathbb{R} \xrightarrow{I_{\text{grad},F}^k} \underline{X}_{\text{grad},F}^k \xrightarrow{G_F^k} \underline{X}_{\text{rot},F}^k \xrightarrow{C_F^k} \mathcal{P}^k(F) \xrightarrow{0} \{0\}$$

Space	V (vertex)	E (edge)	F (polygon)
$\underline{X}_{\text{grad},F}^k$	$\mathbb{R} = \mathcal{P}^k(V)$	$\mathcal{P}^{k-1}(E)$	$\mathcal{P}^{k-1}(F)$
$\underline{X}_{\text{rot},F}^k$		$\mathcal{P}^k(E)$	$\mathcal{R}^{k-1}(F) \times \mathcal{R}^{c,k}(F)$
$\mathcal{P}^k(F)$			$\mathcal{P}^k(F)$

Table: Polynomial components for the two-dimensional spaces

- **Interpolators** = component-wise L^2 -projections
- **Discrete operators** = L^2 -projections of full operator reconstructions

The three-dimensional case I

Exact sequence

$$\mathbb{R} \xrightarrow{I_{\text{grad},T}^k} \underline{X}_{\text{grad},T}^k \xrightarrow{\underline{G}_T^k} \underline{X}_{\text{curl},T}^k \xrightarrow{\underline{C}_T^k} \underline{X}_{\text{div},T}^k \xrightarrow{D_T^k} \mathcal{P}^k(T) \xrightarrow{0} \{0\}$$

Space	V	E	F (face)	T (polyhedron)
$\underline{X}_{\text{grad},T}^k$	\mathbb{R}	$\mathcal{P}^{k-1}(E)$	$\mathcal{P}^{k-1}(F)$	$\mathcal{P}^{k-1}(T)$
$\underline{X}_{\text{curl},T}^k$		$\mathcal{P}^k(E)$	$\mathcal{R}^{k-1}(F) \times \mathcal{R}^{c,k}(F)$	$\mathcal{R}^{k-1}(T) \times \mathcal{R}^{c,k}(T)$
$\underline{X}_{\text{div},T}^k$			$\mathcal{P}^k(F)$	$\mathcal{G}^{k-1}(T) \times \mathcal{G}^{c,k}(T)$
$\mathcal{P}^k(T)$				$\mathcal{P}^k(T)$

Table: Polynomial components for the three-dimensional spaces

Theorem (Exactness of the three-dimensional local DDR sequence)

If the polyhedron T has a trivial topology, this sequence is *exact*.

The three-dimensional case II

Exact sequence

Lemma (Commutative diagram with the sequence of trimmed spaces)

The following commutative diagram holds, expressing the *polynomial consistency* of the discrete vector calculus operators:

$$\begin{array}{ccccccc} \mathcal{P}^{k+1}(T) & \xrightarrow{\text{grad}} & \mathcal{N}^k(T) & \xrightarrow{\text{curl}} & \mathcal{RT}^k(T) & \xrightarrow{\text{div}} & \mathcal{P}^k(T) \\ \downarrow \underline{I}_{\text{grad},T}^k & & \downarrow \underline{I}_{\text{curl},T}^k & & \downarrow \underline{I}_{\text{div},T}^k & & \downarrow i_T \\ \underline{X}_{\text{grad},T}^k & \xrightarrow{\underline{G}_T^k} & \underline{X}_{\text{curl},T}^k & \xrightarrow{\underline{C}_T^k} & \underline{X}_{\text{div},T}^k & \xrightarrow{D_T^k} & \mathcal{P}^k(T) \end{array}$$

The three-dimensional case

Local discrete L^2 -products

- Emulating integration by part formulas, define the **local potentials**

$$\mathbf{P}_{\text{grad},T}^{k+1} : \underline{X}_{\text{grad},T}^k \rightarrow \mathcal{P}^{k+1}(T),$$

$$\mathbf{P}_{\text{curl},T}^k : \underline{X}_{\text{curl},T}^k \rightarrow \mathcal{P}^k(T)^3,$$

$$\mathbf{P}_{\text{div},T}^k : \underline{X}_{\text{div},T}^k \rightarrow \mathcal{P}^k(T)^3$$

- Based on these potentials, we construct **local discrete L^2 -products**

$$(\underline{x}_T, \underline{y}_T)_{\bullet,T} = \underbrace{\int_T P_{\bullet,T} \underline{x}_T \cdot P_{\bullet,T} \underline{y}_T}_{\text{consistency}} + \underbrace{s_{\bullet,T}(\underline{x}_T, \underline{y}_T)}_{\text{stability}} \quad \forall \bullet \in \{\text{grad, curl, div}\}$$

- The L^2 -products are **polynomially exact**

The three-dimensional case

Global sequence

- Let $\Omega \subset \mathbb{R}^3$ as before and let \mathcal{T}_h be a **polyhedral mesh**
- **Global DDR spaces** are defined gluing boundary components:

$$\underline{X}_{\text{grad},h}^k, \quad \underline{X}_{\text{curl},h}^k, \quad \underline{X}_{\text{div},h}^k$$

- **Global operators** are obtained collecting local components:

$$\underline{G}_h^k, \quad \underline{C}_h^k, \quad D_h^k$$

- **Global L^2 -products** $(\cdot, \cdot)_{\bullet,h}$ are obtained assembling element-wise
- The **global DDR sequence** is

$$\mathbb{R} \xrightarrow{I_{\text{grad},h}^k} \underline{X}_{\text{grad},h}^k \xrightarrow{\underline{G}_h^k} \underline{X}_{\text{curl},h}^k \xrightarrow{\underline{C}_h^k} \underline{X}_{\text{div},h}^k \xrightarrow{D_h^k} \mathcal{P}^k(\mathcal{T}_h) \xrightarrow{0} \{0\}$$

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A DDR scheme for magnetostatics

Discrete problem I

- Continuous problem: Find $(\boldsymbol{\sigma}, \mathbf{u}) \in \mathbf{H}(\mathbf{curl}; \Omega) \times \mathbf{H}(\mathbf{div}; \Omega)$ s.t.

$$\int_{\Omega} \boldsymbol{\sigma} \cdot \boldsymbol{\tau} - \int_{\Omega} \mathbf{u} \cdot \mathbf{curl} \boldsymbol{\tau} = 0 \quad \forall \boldsymbol{\tau} \in \mathbf{H}(\mathbf{curl}; \Omega),$$
$$\int_{\Omega} \mathbf{curl} \boldsymbol{\sigma} \cdot \mathbf{v} + \int_{\Omega} \mathbf{div} \mathbf{u} \mathbf{div} \mathbf{v} = \int_{\Omega} \mathbf{f} \cdot \mathbf{v} \quad \forall \mathbf{v} \in \mathbf{H}(\mathbf{div}; \Omega)$$

- The **global bilinear forms** are approximated substituting

$$(\underline{\boldsymbol{\sigma}}_h, \underline{\boldsymbol{\tau}}_h)_{\mathbf{curl}, h} \leftarrow \int_{\Omega} \boldsymbol{\sigma} \cdot \boldsymbol{\tau}$$
$$(\underline{\mathbf{C}}_h^k \underline{\boldsymbol{\tau}}_h, \underline{\mathbf{v}}_h)_{\mathbf{div}, h} \leftarrow \int_{\Omega} \mathbf{curl} \boldsymbol{\tau} \cdot \mathbf{v}$$
$$\int_{\Omega} D_h^k \underline{\mathbf{w}}_h D_h^k \underline{\mathbf{v}}_h \leftarrow \int_{\Omega} \mathbf{div} \mathbf{w} \mathbf{div} \mathbf{v}$$

- The current density linear form is l_h , defined similarly

A DDR scheme for magnetostatics

Discrete problem II

- The **DDR problem** reads: Find $(\underline{\sigma}_h, \underline{u}_h) \in \underline{X}_{\text{curl},h}^k \times \underline{X}_{\text{div},h}^k$ s.t.

$$\begin{aligned}(\underline{\sigma}_h, \underline{\tau}_h)_{\text{curl},h} - (\underline{u}_h, \underline{C}_h^k \underline{\tau}_h)_{\text{div},h} &= 0 & \forall \underline{\tau}_h \in \underline{X}_{\text{curl},h}^k, \\(\underline{C}_h^k \underline{\sigma}_h, \underline{v}_h)_{\text{div},h} + \int_{\Omega} D_h^k \underline{u}_h D_h^k \underline{v}_h &= l_h(\underline{v}_h) & \forall \underline{v}_h \in \underline{X}_{\text{div},h}^k\end{aligned}$$

- **Stability** hinges on the exactness of the portion

$$\mathbb{R} \xrightarrow{I_{\text{grad},h}^k} \underline{X}_{\text{grad},h}^k \xrightarrow{\underline{G}_h^k} \underline{X}_{\text{curl},h}^k \xrightarrow{\underline{C}_h^k} \underline{X}_{\text{div},h}^k \xrightarrow{D_h^k} \mathcal{P}^k(\mathcal{T}_h) \xrightarrow{0} \{0\}$$

A DDR scheme for magnetostatics

Global exactness I

Theorem (Exactness properties of the global DDR sequence)

Let $\Omega \subset \mathbb{R}^3$ be an open connected polyhedral domain. Then, it holds

$$\text{Im } D_h^k = \mathcal{P}^k(\mathcal{T}_h).$$

If Ω *does not enclose any void*, we additionally have

$$\text{Im } \underline{C}_h^k = \text{Ker } D_h^k.$$

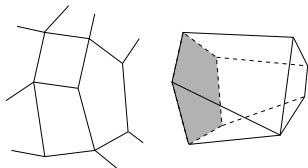
- $\text{Im } D_h^k = \mathcal{P}^k(\mathcal{T}_h)$ follows from the classical Fortin's argument
- The inclusion $\text{Im } \underline{C}_h^k \subset \text{Ker } D_h^k$ results from **local exactness**
- We prove $\text{Ker } D_h^k \subset \text{Im } \underline{C}_h^k$ in two steps. Let $\underline{v}_h \in \text{Ker } D_h^k$. Then:
 - **Local exactness** gives $\underline{\tau}_T \in \underline{X}_{\text{curl},T}^k$ s.t. $\underline{v}_T = \underline{C}_T^k \underline{\tau}_T$ for all $T \in \mathcal{T}_h$
 - The local vectors are then **glued together**

A DDR scheme for magnetostatics

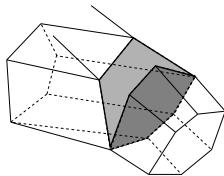
Global exactness II

To glue together local vectors, we use the fact that the mesh can be topologically assembled by **a succession of the following operations**:

- 1 Add a new element by glueing one ot its faces to an element in the mesh



- 2 Glue together two faces of elements in the mesh s.t. the edges along which the faces are already glued together **form a connected path**



This is only possible since Ω does not enclose any void!

A DDR scheme for magnetostatics

Stability and well-posedness

Theorem (Well-posedness)

Let $\Omega \subset \mathbb{R}^3$ be an open *simply connected* polyhedral domain *that does not enclose any void*. Then, $(\underline{\sigma}_h, \underline{u}_h) \in \underline{X}_{\text{curl},h}^k \times \underline{X}_{\text{div},h}^k$ is unique and there exists $C > 0$ independent of h s.t.

$$\|\underline{\sigma}_h\|_{\text{curl},h} + \|\underline{C}_h^k \underline{\sigma}_h\|_{\text{div},h} + \|\underline{u}_h\|_{\text{div},h} + \|D_h^k \underline{u}_h\|_{L^2(\Omega)} \leq C \|f\|_{\Omega}.$$

Proof.

Analogous to the continuous case since all the relevant properties have been reproduced at the discrete level. □

Numerical examples

Setting

- Let $(\mathcal{T}_h)_{h \in \mathcal{H}}$ be a **regular polyhedral mesh sequence**
- We consider a known solution $(\boldsymbol{\sigma}, \mathbf{u})$ to assess **convergence rate r** s.t.

approximation error $\propto h^r$

- The error

$$(\underline{\boldsymbol{e}}_h, \underline{\boldsymbol{\varepsilon}}_h) := (\underline{\boldsymbol{\sigma}}_h - \underline{\mathbf{I}}_{\text{curl},h}^k \boldsymbol{\sigma}, \underline{\mathbf{u}}_h - \underline{\mathbf{I}}_{\text{div},h}^k \mathbf{u})$$

is measured in the natural **energy norm** s.t.

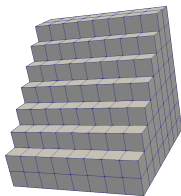
$$\|(\underline{\boldsymbol{e}}_h, \underline{\boldsymbol{\varepsilon}}_h)\|_{\text{en},h} := [(\underline{\boldsymbol{e}}_h, \underline{\boldsymbol{e}}_h)_{\text{curl},h} + (\underline{\boldsymbol{\varepsilon}}_h, \underline{\boldsymbol{\varepsilon}}_h)_{\text{div},h}]^{\frac{1}{2}}$$

- The implementation is based on the **HArDCore3D** C++ library¹

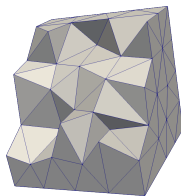
¹See <https://tinyurl.com/HarDCore3D>

Numerical examples

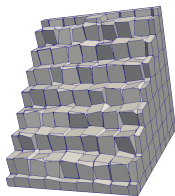
Meshes



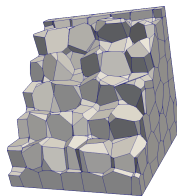
(a) Cubic-Cells



(b) Tetgen-Cube-0



(c) Voro-small-0



(d) Voro-small-1

Figure: Mesh families used in the numerical tests

Numerical examples

Convergence in the energy norm

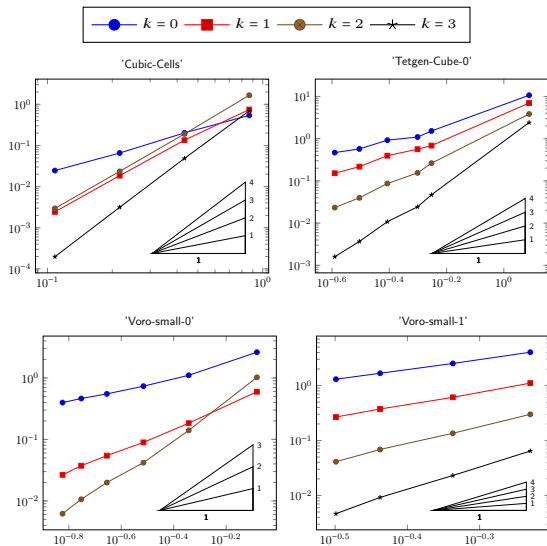


Figure: Energy error versus mesh size h . We have $\|(\underline{e}_h, \underline{\varepsilon}_h)\|_{en,h} \propto h^{k+1}$

Conclusions and perspectives

- A **novel approach** for the numerical solution of PDE problems
- **Key features:** support of general polyhedral meshes and high-order
- **Novel computational strategies** made possible
- Natural extensions to **variable coefficients** and **nonlinearities**

- **Applications** (electromagnetism, incompressible fluid mechanics, . . .)
- Formalization using **differential forms** (ongoing)
- Development of **novel sequences** (e.g., elasticity)
- . . .

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