Hybrid High-Order methods for nonlinear problems

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Two crucial problems for humanity





Hybrid High-Order (HHO) methods



Figure: Examples of supported meshes $\mathcal{M}_h = (\mathcal{T}_h, \mathcal{F}_h)$ in 2d and 3d

- Capability of handling general polyhedral meshes
- Construction valid for arbitrary space dimensions
- Arbitrary approximation order (including k = 0)
- Natural extension to nonlinear problems
- Reduced computational cost after static condensation
- Key idea: replace spaces and operators with discrete counterparts

References for this presentation

HHO for Leray–Lions problems

- Analysis tools and convergence [DP and Droniou, 2017a]
- Basic error estimates [DP and Droniou, 2017b]
- Stabilization-free [DP, Droniou, Manzini, 2018]
- Improved estimates (general meshes) [DP, Droniou, Harnist, 2021]
- Improved estimates (standard meshes) [Carstensen and Tran, 2020]
- Applications
 - Nonlinear elasticity [Botti, DP, Sochala, 2017]
 - Nonlinear poroelasticity [Botti, DP, Sochala, 2018]
 - Non-Newtonian fluids [Botti, Castanon Quiroz, DP, Harnist, 2020]
- General introduction to HHO methods:

Di Pietro, D. A. and Droniou, J. (2020). **The Hybrid High-Order method for polytopal meshes. Design, analysis, and applications**, volume 19 of *Modeling, Simulation and Application*. Springer International Publishing.

Outline

1 Leray–Lions problems

2 Creeping flows of non-Newtonian fluids

Model problem

- \blacksquare Let $\Omega \subset \mathbb{R}^d$ denote a bounded connected polyhedral domain
- Let $r \in (0, +\infty)$ and $r' \coloneqq \frac{r}{r-1}$
- Consider the problem: Given $f \in L^{r'}(\Omega)$, find $u : \Omega \to \mathbb{R}$ s.t.

$$-\nabla \cdot \boldsymbol{\sigma}(\boldsymbol{x}, \nabla \boldsymbol{u}) = f \quad \text{in } \Omega,$$
$$\boldsymbol{u} = 0 \quad \text{on } \partial \Omega$$

In weak formulation: Find $u \in W_0^{1,r}(\Omega)$ s.t.

$$\int_{\Omega} \boldsymbol{\sigma}(\cdot, \boldsymbol{\nabla} \boldsymbol{u}) \cdot \boldsymbol{\nabla} \boldsymbol{v} = \int_{\Omega} f \boldsymbol{v} \qquad \forall \boldsymbol{v} \in W_0^{1,r}(\Omega).$$

The key differential operator is the gradient

Flux function

Assumption (Flux function I)

The Carathéodory function¹ $\sigma : \Omega \times \mathbb{R}^d \to \mathbb{R}^d$ is s.t., for a.e. $x \in \Omega$ and all $\eta, \xi \in \mathbb{R}^d$,

• Growth. There exists a real number $\overline{\sigma} > 0$ s.t.

 $|\sigma(x,\eta) - \sigma(x,0)| \le \overline{\sigma}|\eta|^{r-1}.$

• Coercivity. There is a real number $\underline{\sigma} > 0$ s.t.,

 $\sigma(\boldsymbol{x},\boldsymbol{\eta})\cdot\boldsymbol{\eta}\geq\underline{\sigma}|\boldsymbol{\eta}|^r.$

Monotonicity. It holds

$$(\boldsymbol{\sigma}(\boldsymbol{x},\boldsymbol{\eta}) - \boldsymbol{\sigma}(\boldsymbol{x},\boldsymbol{\xi})) \cdot (\boldsymbol{\eta} - \boldsymbol{\xi}) \ge 0.$$

 ${}^{1}\sigma(x,\cdot)$ continuous, $\sigma(\cdot,\eta)$ measurable

L^2 -orthogonal projectors on local polynomial spaces

- Let a polynomial degree $k \ge 0$ and a mesh element or face X be fixed
- Define the polynomial space

 $\mathbb{P}^k(X) \coloneqq \{$ restriction to X of d-variate polynomials of total degree $\leq k \}$

• The L^2 -orthogonal projector $\pi^k_X : L^2(X) \to \mathbb{P}^k(X)$ is s.t.

$$\int_X (\pi_X^k v - v) w = 0 \text{ for all } w \in \mathbb{P}^k(X)$$

• Optimal approximation properties hold [DP and Droniou, 2020]

A key remark

- Let a polytopal mesh element $T \in \mathcal{T}_h$ be fixed
- Recall the following IBP formula, valid for all $(v, \tau) \in W^{1,1}(T) \times C^{\infty}(\overline{T})^d$:

$$\int_{T} \boldsymbol{\nabla} \boldsymbol{v} \cdot \boldsymbol{\tau} = -\int_{T} \boldsymbol{v} \ (\boldsymbol{\nabla} \cdot \boldsymbol{\tau}) + \sum_{F \in \mathcal{F}_{T}} \int_{F} \boldsymbol{v} \ (\boldsymbol{\tau} \cdot \boldsymbol{n}_{TF})$$

 \blacksquare Given an integer $k \geq 0,$ taking $\pmb{\tau} \in \mathbb{P}^k(T)^d$ we can write

$$\int_{T} \pi_{T}^{k} (\nabla v) \cdot \tau = - \int_{T} \pi_{T}^{k} v (\nabla \cdot \tau) + \sum_{F \in \mathcal{F}_{T}} \int_{F} \pi_{F}^{k} v_{|F|} (\tau \cdot \mathbf{n}_{TF})$$

• Hence, $\pi_T^k(\nabla v)$ can be computed from $\pi_T^k v$ and $(\pi_F^k v|_F)_{F \in \mathcal{F}_T}$!

Local HHO space and interpolator



Figure: \underline{U}_T^k for $k \in \{0, 1, 2\}$ and d = 2

- For $k \ge 0$ and $T \in \mathcal{T}_h$, define the local HHO space
 - $\underline{U}_{T}^{k} \coloneqq \left\{ \underline{v}_{T} = (v_{T}, (v_{F})_{F \in \mathcal{F}_{T}}) : v_{T} \in \mathbb{P}^{k}(T) \text{ and } v_{F} \in \mathbb{P}^{k}(F) \text{ for all } F \in \mathcal{F}_{T} \right\}$
- The local interpolator $\underline{I}_T^k: W^{1,1}(T) \to \underline{U}_T^k$ is s.t., for all $v \in W^{1,1}(T)$,

$$\underline{I}_T^k v \coloneqq \left(\pi_T^k v, (\pi_F^k v)_{F \in \mathcal{F}_T} \right)$$

Gradient reconstruction

• Let $T \in \mathcal{T}_h$. We define the local gradient reconstruction

$$G_T^k : \underline{U}_T^k \to \mathbb{P}^k(T)^d$$

s.t., for all
$$\underline{v}_T \in \underline{U}_T^k$$
,

$$\int_T \mathbf{G}_T^k \underline{v}_T \cdot \boldsymbol{\tau} = -\int_T v_T \ (\boldsymbol{\nabla} \cdot \boldsymbol{\tau}) + \sum_{F \in \mathcal{F}_T} \int_F v_F \ (\boldsymbol{\tau} \cdot \boldsymbol{n}_{TF}) \quad \forall \boldsymbol{\tau} \in \mathbb{P}^k(T)^d$$

By construction, we have,

$$\boldsymbol{G}_T^k(\underline{I}_T^k\boldsymbol{v}) = \boldsymbol{\pi}_T^k(\boldsymbol{\nabla}\boldsymbol{v}) \quad \forall \boldsymbol{v} \in W^{1,1}(T)$$

• $(G_T^k \circ \underline{I}_T^k)$ therefore has optimal approximation properties in $\mathbb{P}^k(T)^d$

Global HHO space and gradient reconstruction

■ The global HHO space is obtained patching interface unknowns:

$$\begin{split} \underline{U}_{h}^{k} \coloneqq \left\{ \underline{v}_{h} = ((v_{T})_{T \in \mathcal{T}_{h}}, (v_{F})_{F \in \mathcal{T}_{h}}) : \\ v_{T} \in \mathbb{P}^{k}(T) \text{ for all } T \in \mathcal{T}_{h} \text{ and } v_{F} \in \mathbb{P}^{k}(F) \text{ for all } F \in \mathcal{T}_{h} \end{split} \right\}$$

• The global gradient
$$G_h^k : \underline{U}_h^k \to \mathbb{P}^k(\mathcal{T}_h)^d$$
 is s.t.

$$\forall \underline{v}_h \in \underline{U}_h^k, \quad (\boldsymbol{G}_h^k \underline{v}_h)_{|T} \coloneqq \boldsymbol{G}_T^k \underline{v}_T \quad \forall T \in \mathcal{T}_h$$

Accounting for boundary conditions, we set

$$\underline{U}_{h,0}^{k} \coloneqq \left\{ \underline{v}_{h} \in \underline{U}_{h}^{k} : v_{F} = 0 \text{ for all } F \in \mathcal{F}_{h} \text{ s.t. } F \subset \partial \Omega \right\}$$

Discrete Sobolev norms

- We need to endow \underline{U}_h^k with a Sobolev structure
- We define the discrete Sobolev seminorm s.t., for all $\underline{v}_h \in \underline{U}_h^k$,

$$\|\underline{v}_{h}\|_{1,r,h} \coloneqq \left(\sum_{T \in \mathcal{T}_{h}} \|\underline{v}_{T}\|_{1,r,T}^{r}\right)^{\frac{1}{r}}$$

where, for all $T \in \mathcal{T}_h$,

$$\|\underline{\boldsymbol{v}}_{T}\|_{1,r,T} \coloneqq \left(\|\boldsymbol{\nabla}\boldsymbol{v}_{T}\|_{L^{r}(T)^{d}}^{r} + \sum_{F \in \mathcal{F}_{T}} h_{F}^{1-r} \|\boldsymbol{v}_{F} - \boldsymbol{v}_{T}\|_{L^{r}(F)}^{r} \right)^{\frac{1}{r}}$$

• The factor h_F^{1-r} in the boundary term ensures the appropriate scaling

Discrete functional analysis results I

Theorem (Discrete Sobolev–Poincaré inequalities)

Let

$$1 \le q \le \frac{dr}{d-r}$$
 if $1 \le r < d$ and $1 \le q < +\infty$ if $r \ge d$.

Then, for all $\underline{v}_h \in \underline{U}_{h,0}^k$, letting $v_h \in \mathbb{P}^k(\mathcal{T}_h)$ be s.t.

$$(v_h)_{|T} \coloneqq v_T \qquad \forall T \in \mathcal{T}_h,$$

it holds, with C > 0 depending only on Ω , k, r, q, and mesh regularity,

$$\|v_h\|_{L^q(\Omega)} \le C \|\underline{v}_h\|_{1,r,h}.$$

Corollary (Discrete Sobolev norms)

The mapping $\|\cdot\|_{1,r,h}$ is a norm on $\underline{U}_{h,0}^k$.

Discrete functional analysis results II

Theorem (Discrete compactness)

Let $(\mathcal{M}_h)_{h>0}$ be a regular mesh sequence and $(\underline{v}_h)_{h>0} \in (\underline{U}_{h,0}^k)_{h>0}$ s.t.

$$\|\underline{v}_h\|_{1,r,h} \leq C$$
 for all $h > 0$.

Then, there exists $v \in W_0^{1,r}(\Omega)$ s.t., up to a subsequence as $h \to 0$,

• $v_h \rightarrow v$ strongly in $L^q(\Omega)$ for all $1 \le q < \begin{cases} \frac{dr}{d-r} & \text{if } r < d, \\ +\infty & \text{otherwise;} \end{cases}$

•
$$G_{h\underline{\nu}_{h}}^{k} \rightarrow \nabla v$$
 weakly in $L^{r}(\Omega)^{d}$.

Proposition (Strong convergence of the gradient for smooth functions)

With $(\mathcal{M}_h)_{h>0}$ as before it holds, for all $\varphi \in W^{1,r}(\Omega)$,

$$G_h^k(\underline{I}_h^k\varphi) \to \nabla \varphi$$
 strongly in $L^r(\Omega)^d$ as $h \to 0$.

• Define the function $a_h : \underline{U}_h^k \times \underline{U}_h^k \to \mathbb{R}$ s.t.

$$\mathbf{a}_h(\underline{w}_h,\underline{v}_h)\coloneqq \int_{\Omega} \boldsymbol{\sigma}(\cdot,\boldsymbol{G}_h^k\underline{w}_h)\cdot\boldsymbol{G}_h^k\underline{v}_h + \sum_{T\in\mathcal{T}_h}\mathbf{s}_T(\underline{w}_T,\underline{v}_T)$$

• Above, s_T is a stabilization obtained penalizing face residuals s.t.

- $\blacksquare \|G_T^k \underline{v}_T\|_{L^r(T)^d}^r + \mathbf{s}_T(\underline{v}_T, \underline{v}_T) \simeq \|\underline{v}_T\|_{1,r,T}^r \text{ uniformly in } h$
- $s_T(\underline{I}_T^k w, \underline{v}_T) = 0$ for all $(w, \underline{v}_T) \in \mathbb{P}^{k+1}(T) \times \underline{U}_T^k$
- Hölder continuity and strong monotonicity hold

Discrete problem II

The discrete Leray-Lions problem reads:

Find
$$\underline{u}_h \in \underline{U}_{h,0}^k$$
 s.t. $a_h(\underline{u}_h, \underline{v}_h) = \int_{\Omega} f v_h \quad \forall \underline{v}_h \in \underline{U}_{h,0}^k$

Lemma (Existence and a priori bound)

There is at least one solution $\underline{u}_h \in \underline{U}_{h,0}^k$, and any solution satisfies

$$\|\underline{u}_{h}\|_{1,r,h} \leq C \|f\|_{L^{r'}(\Omega)}^{\frac{1}{r-1}},$$

with real number C > 0 independent of h.

Remark (Uniqueness)

Uniqueness holds replacing monotonicity with strict monotonicity.

Theorem (Convergence)

Let $(\mathcal{M}_h)_{h>0}$ be a regular mesh sequence and let $(\underline{u}_h)_{h>0}$ be the corresponding sequence of discrete solutions. Then, as $h \to 0$, up to a subsequence,

•
$$u_h \to u$$
 strongly in $L^q(\Omega)$ with $1 \le q < \begin{cases} \frac{dr}{d-r} & \text{if } r < d, \\ +\infty & \text{otherwise,} \end{cases}$

• $G_{h}^{k}\underline{u}_{h} \rightarrow \nabla u$ weakly in $L^{r}(\Omega)^{d}$,

with $u \in W_0^{1,r}(\Omega)$ solution to the continuous problem. If, additionally, σ is strictly monotone, then u is unique and $G_h^k \underline{u}_h$ converges strongly.

Proof.

- Combining the a priori bound with discrete compactness, we infer the existence of $u \in W_0^{1,r}(\Omega)$ s.t. the above convergences hold
- Taking $\underline{v}_h = \underline{I}_h^k \varphi$ as test function with $\varphi \in C_c^{\infty}(\Omega)$ and using Minty's trick, we infer that u solves the continuous problem
- Using Vitali's theorem, we prove strong convergence of $G_h^k \underline{u}_h$ under strict monotonicity of σ

Error estimates I

Assumption (Flux function II)

In addition to Assumption I, it holds, for a.e. $x \in \Omega$ and all $\eta, \xi \in \mathbb{R}^d$,

Hölder continuity. There exists a real number $\sigma^* > 0$ s.t.

$$|\boldsymbol{\sigma}(\boldsymbol{x},\boldsymbol{\eta}) - \boldsymbol{\sigma}(\boldsymbol{x},\boldsymbol{\xi})| \leq \sigma^* |\boldsymbol{\eta} - \boldsymbol{\xi}| \left(|\boldsymbol{\eta}|^{r-2} + |\boldsymbol{\xi}|^{r-2} \right).$$

Strong monotonicity. There exists a real number $\sigma_* > 0$ s.t.

$$(\boldsymbol{\sigma}(\boldsymbol{x},\boldsymbol{\eta}) - \boldsymbol{\sigma}(\boldsymbol{x},\boldsymbol{\xi})) \cdot (\boldsymbol{\eta} - \boldsymbol{\xi}) \geq \sigma_* |\boldsymbol{\eta} - \boldsymbol{\xi}|^2 (|\boldsymbol{\eta}| + |\boldsymbol{\xi}|)^{r-2}.$$

Remark (*r*-Laplacian)

The above assumptions are verified by the r-Laplace flux function

$$\boldsymbol{\sigma}(\boldsymbol{x},\boldsymbol{\eta}) = |\boldsymbol{\eta}|^{r-2}\boldsymbol{\eta}.$$

Error estimates II

Theorem (Basic error estimate)

Assume $u \in W^{k+2,r}(\mathcal{T}_h)$ and $\sigma(\cdot, \nabla u) \in W^{k+1,r'}(\mathcal{T}_h)^d$ and let • if $r \ge 2$,

$$\mathcal{E}_{h}(u) \coloneqq h^{k+1} |u|_{W^{k+2,r}(\mathcal{T}_{h})} + \frac{h^{\frac{k+1}{r-1}}}{h^{\frac{k+1}{r-1}}} \left(|u|_{W^{k+2,r}(\mathcal{T}_{h})}^{\frac{1}{r-1}} + |\sigma(\cdot, \nabla u)|_{W^{k+1,r'}(\mathcal{T}_{h})^{d}}^{\frac{1}{r-1}} \right);$$

■ *if* r < 2,

$$\mathcal{E}_h(u) \coloneqq h^{(k+1)(r-1)} |u|_{W^{k+2,r}(\mathcal{T}_h)}^{r-1} + h^{k+1} |\sigma(\cdot, \nabla u)|_{W^{k+1,r'}(\mathcal{T}_h)^d}.$$

Then, it holds

$$|\underline{I}_{h}^{k}u - \underline{u}_{h}\|_{1,r,h} \leq C\mathcal{E}_{h}(u),$$

with C > 0 depending only on Ω , k, r, σ , $\overline{\sigma}$, σ_* , σ^* , and mesh regularity.

Improved error estimates

The above estimate gives the following asymptotic convergence rates:

$$\begin{cases} h^{\frac{k+1}{r-1}} & \text{if } r \ge 2, \\ h^{(k+1)(r-1)} & \text{if } 1 < r < 2 \end{cases}$$

Successively [DP, Droniou, Harnist, 2021] proved

 h^{k+1} in the non-degenerate case for $1 < r \le 2$,

with intermediate rates depending on a degeneracy parameterVery recently, [Carstensen and Tran, 2020] proved convergence in

$$h^{\frac{k+1}{3-r}}$$
 for $1 < r \le 2$

for a variation of the HHO method on conforming simplicial meshes based on a stable gradient inspired by [DP, Droniou, Manzini, 2018]

Numerical example

Convergence for r = 3

h	$\ \underline{I}_h^k u - \underline{u}_h\ _{1,r,h}$	EOC	h	$\ \underline{I}_h^k u - \underline{u}_h\ _{1,r,h}$	EOC
	k = 1 (1)			k = 1 (1)	
$3.07 \cdot 10^{-2}$	$1.71 \cdot 10^{-2}$	_	$6.5 \cdot 10^{-2}$	$3.06 \cdot 10^{-2}$	_
$1.54 \cdot 10^{-2}$	$4.72 \cdot 10^{-3}$	1.87	$3.15 \cdot 10^{-2}$	$1.1 \cdot 10^{-2}$	1.41
$7.68 \cdot 10^{-3}$	$1.16 \cdot 10^{-3}$	2.02	$1.61 \cdot 10^{-2}$	$3.35 \cdot 10^{-3}$	1.77
$3.84 \cdot 10^{-3}$	$2.96 \cdot 10^{-4}$	1.97	$9.09 \cdot 10^{-3}$	$1.25 \cdot 10^{-3}$	1.72
$1.92\cdot 10^{-3}$	$7.77\cdot 10^{-5}$	1.93	$4.26\cdot 10^{-3}$	$3.58\cdot10^{-4}$	1.65
	$k=2 \left(\frac{3}{2}\right)$			$k=2 \left(\frac{3}{2}\right)$	
$3.07\cdot 10^{-2}$	$2.72\cdot 10^{-3}$	_	$6.5 \cdot 10^{-2}$	$1.18\cdot 10^{-2}$	_
$1.54 \cdot 10^{-2}$	$2.32 \cdot 10^{-4}$	3.57	$3.15 \cdot 10^{-2}$	$2.33 \cdot 10^{-3}$	2.24
$7.68 \cdot 10^{-3}$	$3.32 \cdot 10^{-5}$	2.79	$1.61 \cdot 10^{-2}$	$4.4 \cdot 10^{-4}$	2.48
$3.84 \cdot 10^{-3}$	$7.25 \cdot 10^{-6}$	2.2	$9.09 \cdot 10^{-3}$	$1.02 \cdot 10^{-4}$	2.56
$1.92\cdot 10^{-3}$	$1.81\cdot 10^{-6}$	2.00	$4.26\cdot 10^{-3}$	$1.42\cdot 10^{-5}$	2.60
	k = 3 (2)			k = 3 (2)	
$3.07\cdot 10^{-2}$	$3.1\cdot 10^{-4}$	_	$6.5 \cdot 10^{-2}$	$2.75\cdot 10^{-3}$	_
$1.54 \cdot 10^{-2}$	$2.97 \cdot 10^{-5}$	3.4	$3.15 \cdot 10^{-2}$	$2.69 \cdot 10^{-4}$	3.21
$7.68 \cdot 10^{-3}$	$4.4 \cdot 10^{-6}$	2.74	$1.61 \cdot 10^{-2}$	$4.01 \cdot 10^{-5}$	2.84
$3.84 \cdot 10^{-3}$	$9.76 \cdot 10^{-7}$	2.17	$9.09 \cdot 10^{-3}$	$1.31 \cdot 10^{-5}$	1.96
$1.92 \cdot 10^{-3}$	$2.41 \cdot 10^{-7}$	2.02	$4.26 \cdot 10^{-3}$	$2.21 \cdot 10^{-6}$	2.35

Table: Triangular mesh family

Table: Voronoi mesh family

Outline



2 Creeping flows of non-Newtonian fluids

Model problem I

• Let $d \in \{2,3\}$. Given $f : \Omega \to \mathbb{R}^d$, the nonlinear Stokes problem reads: Find $u : \Omega \to \mathbb{R}^d$ and $p : \Omega \to \mathbb{R}$ s.t.

$$\begin{aligned} - \nabla \cdot \boldsymbol{\sigma} (\boldsymbol{\nabla}_{\mathrm{s}} \boldsymbol{u}) + \boldsymbol{\nabla} p &= \boldsymbol{f} & \text{ in } \boldsymbol{\Omega}, \\ \boldsymbol{\nabla} \cdot \boldsymbol{u} &= \boldsymbol{0} & \text{ in } \boldsymbol{\Omega}, \\ \boldsymbol{u} &= \boldsymbol{0} & \text{ on } \partial \boldsymbol{\Omega}, \\ \int_{\boldsymbol{\Omega}} p &= \boldsymbol{0}, \end{aligned}$$

• We focus, for the sake of simplicity, on power-law fluids, for which

$$\boldsymbol{\sigma}(\boldsymbol{\tau}) = |\boldsymbol{\tau}|^{r-2} \boldsymbol{\tau} \quad \forall \boldsymbol{\tau} \in \mathbb{R}^{d \times d}_{\mathrm{sym}}$$

For r ∈ (1,2] the fluid is shear-thinning, for r ≥ 2, shear-thickening
More general strain rate-shear stress laws can be considered

Model problem II

Define the following spaces:

$$\boldsymbol{U} \coloneqq W_0^{1,r}(\Omega)^d, \quad \boldsymbol{P} \coloneqq \left\{ \boldsymbol{q} \in \boldsymbol{L}^{r'}(\Omega) \ : \ \int_{\Omega} \boldsymbol{q} = \boldsymbol{0} \right\}$$

Taking $f \in L^{r'}(\Omega)^d$, the weak formulation is: Find $(u, p) \in U \times P$ s.t.

$$a(\boldsymbol{u}, \boldsymbol{v}) + b(\boldsymbol{v}, p) = \int_{\Omega} \boldsymbol{f} \cdot \boldsymbol{v} \quad \forall \boldsymbol{v} \in \boldsymbol{U},$$
$$-b(\boldsymbol{u}, q) = 0 \qquad \forall p \in \boldsymbol{P}$$

where $a: U \times U \to \mathbb{R}$ and $b: U \times P \to \mathbb{R}$ are s.t.

$$a(\mathbf{w},\mathbf{v})\coloneqq\int_{\Omega}\boldsymbol{\sigma}(\boldsymbol{\nabla}_{\mathrm{s}}\mathbf{w}):\boldsymbol{\nabla}_{\mathrm{s}}\mathbf{v},\quad b(\mathbf{v},q)\coloneqq-\int_{\Omega}(\boldsymbol{\nabla}\cdot\mathbf{v})\ q$$

The extension of stability results is non-trivial

Given $T \in \mathcal{T}_h$, the vector version of the local HHO space is

$$\underline{U}_{T}^{k} \coloneqq \left\{ \underline{\nu}_{T} = (\nu_{T}, (\nu_{F})_{F \in \mathcal{F}_{T}}) : \nu_{T} \in \mathbb{P}^{k}(T)^{d} \text{ and } \nu_{F} \in \mathbb{P}^{k}(F)^{d} \text{ for all } F \in \mathcal{F}_{T} \right\}$$

• We furnish \underline{U}_T^k with the strain rate $W^{1,r}$ -like seminorm

$$\|\underline{\boldsymbol{\nu}}_T\|_{\varepsilon,r,T} \coloneqq \left(\|\boldsymbol{\nabla}_{\mathbf{s}} \boldsymbol{\nu}_T\|_{L^r(T)^{d \times d}}^r + \sum_{F \in \mathcal{F}_T} h_F^{r-1} \|\boldsymbol{\nu}_F - \boldsymbol{\nu}_T\|_{L^r(F)^d}^r \right)^{\frac{1}{r}}$$

Notice that the symmetric gradient replaces the gradient!

Symmetric gradient and divergence reconstructions

The local symmetric gradient reconstruction is s.t.

$$\boldsymbol{G}_{\mathrm{s},T}^{k}: \underline{\boldsymbol{U}}_{T}^{k} \to \mathbb{P}^{k}(T; \mathbb{R}_{\mathrm{sym}}^{d \times d})$$

s.t., for all $\underline{v}_T \in \underline{U}_T^k$ and all $\tau \in \mathbb{P}^k(T; \mathbb{R}_{\mathrm{sym}}^{d \times d})$,

$$\int_T \boldsymbol{G}_{\mathrm{s},T}^k \underline{\boldsymbol{\nu}_T} : \boldsymbol{\tau} = -\int_T \boldsymbol{\nu_T} \cdot (\boldsymbol{\nabla} \cdot \boldsymbol{\tau}) + \sum_{F \in \mathcal{F}_T} \int_F \boldsymbol{\nu_F} \cdot (\boldsymbol{\tau} \boldsymbol{n}_{TF})$$

• A divergence reconstruction $D_T^k : \underline{U}_T^k \to \mathbb{P}^k(T)$ is obtained setting

$$D_T^k \coloneqq \operatorname{tr}(\boldsymbol{G}_{\mathrm{s},T}^k)$$

• With \underline{I}_T^k interpolator on \underline{U}_T^k we have, for all $v \in W^{1,1}(T)^d$,

$$\boldsymbol{G}_{\mathrm{s},T}^{k}(\boldsymbol{\underline{I}}_{T}^{k}\boldsymbol{\nu}) = \boldsymbol{\pi}_{T}^{k}(\boldsymbol{\nabla}_{\mathrm{s}}\boldsymbol{\nu}), \qquad \boldsymbol{D}_{T}^{k}(\boldsymbol{\underline{I}}_{T}^{k}\boldsymbol{\nu}) = \boldsymbol{\pi}_{T}^{k}(\boldsymbol{\nabla}\boldsymbol{\cdot}\boldsymbol{\nu})$$

Global HHO space and strain reconstruction

At the global level, we define the velocity space

$$\begin{split} \underline{U}_{h}^{k} &\coloneqq \left\{ \underline{v}_{h} = ((v_{T})_{T \in \mathcal{T}_{h}}, (v_{F})_{F \in \mathcal{T}_{h}}) : \\ v_{T} \in \mathbb{P}^{k}(T)^{d} \text{ for all } T \in \mathcal{T}_{h} \text{ and } v_{F} \in \mathbb{P}^{k}(F)^{d} \text{ for all } F \in \mathcal{F}_{h} \right\} \end{split}$$

along with its subspace with strongly enforced BC

$$\underline{U}_{h,0}^k \coloneqq \left\{ \underline{\nu}_h \in \underline{U}_h^k \, : \, \mathbf{\nu}_F = \mathbf{0} \text{ for all } F \in \mathcal{F}_h \text{ s.t. } F \subset \partial \Omega \right\}$$

- We furnish $\underline{U}_{h,0}^k$ with the global $W^{1,r}$ -seminorm $\|\cdot\|_{\varepsilon,r,h}$
- The global strain reconstruction $G_{s,h}^k : \underline{U}_h^k \to \mathbb{P}^k(\mathcal{T}_h; \mathbb{R}_{sym}^{d \times d})$ is s.t.

$$\forall \underline{\mathbf{v}}_h \in \underline{\mathbf{U}}_h^k, \quad (\mathbf{G}_{\mathrm{s},h}^k \underline{\mathbf{v}}_h)_{|T} \coloneqq \mathbf{G}_{\mathrm{s},T}^k \underline{\mathbf{v}}_T \quad \forall T \in \mathcal{T}_h$$

• The viscous function $a_h : \underline{U}_h^k \times \underline{U}_h^k \to \mathbb{R}$ is s.t.

$$\mathbf{a}_h(\underline{\boldsymbol{u}}_h,\underline{\boldsymbol{v}}_h) \coloneqq \int_{\Omega} \boldsymbol{\sigma}(\boldsymbol{G}_{\mathrm{s},h}^k \underline{\boldsymbol{u}}_h) : \boldsymbol{G}_{\mathrm{s},h}^k \underline{\boldsymbol{v}}_h + \sum_{T \in \mathcal{T}_h} \mathbf{s}_T(\underline{\boldsymbol{u}}_T,\underline{\boldsymbol{v}}_T)$$

• To formulate assumptions on s_T , we introduce the singular exponent

 $\tilde{r} \coloneqq \min(r, 2)$

Viscous function II

Assumption

The stabilization function s_T is linear in its second argument and it satisfies:

- Stability. For all $\underline{\nu}_T \in \underline{U}_T^k$, $\|G_{s,T}^k \underline{\nu}_T\|_{L^r(T)^{d \times d}}^2 + s_T(\underline{\nu}_T, \underline{\nu}_T) \simeq \|\underline{\nu}_T\|_{\varepsilon,r,T}^2$
- Polynomial consistency. For all $(w, \underline{v}_T) \in \mathbb{P}^{k+1}(T)^d \times \underline{U}_T^k$, $s_T(\underline{I}_T^k w, \underline{v}_T) = 0$
- **Hölder continuity.** For all $\underline{u}_T, \underline{v}_T, \underline{w}_T \in \underline{U}_T^k$, setting $\underline{e}_T \coloneqq \underline{u}_T \underline{w}_T$,

$$\begin{split} |\mathbf{s}_{T}(\underline{\boldsymbol{u}}_{T},\underline{\boldsymbol{v}}_{T}) - \mathbf{s}_{T}(\underline{\boldsymbol{w}}_{T},\underline{\boldsymbol{v}}_{T})| \lesssim \\ & \left(\mathbf{s}_{T}(\underline{\boldsymbol{u}}_{T},\underline{\boldsymbol{u}}_{T}) + \mathbf{s}_{T}(\underline{\boldsymbol{w}}_{T},\underline{\boldsymbol{w}}_{T})\right)^{\frac{r-\bar{r}}{r}} \mathbf{s}_{T}(\underline{\boldsymbol{e}}_{T},\underline{\boldsymbol{e}}_{T})^{\frac{\bar{r}-1}{r}} \mathbf{s}_{T}(\underline{\boldsymbol{v}}_{T},\underline{\boldsymbol{v}}_{T})^{\frac{1}{r}} \end{split}$$

Strong monotonicity. For all $\underline{u}_T, \underline{w}_T \in \underline{U}_T^k$, setting $\underline{e}_T \coloneqq \underline{u}_T - \underline{w}_T$,

$$\left(\mathrm{s}_{T}\left(\underline{\boldsymbol{u}}_{T},\underline{\boldsymbol{e}}_{T}\right)-\mathrm{s}_{T}\left(\underline{\boldsymbol{w}}_{T},\underline{\boldsymbol{e}}_{T}\right)\right)\left(\mathrm{s}_{T}\left(\underline{\boldsymbol{u}}_{T},\underline{\boldsymbol{u}}_{T}\right)+\mathrm{s}_{T}\left(\underline{\boldsymbol{w}}_{T},\underline{\boldsymbol{w}}_{T}\right)\right)^{\frac{2-\bar{r}}{r}}\gtrsim \mathrm{s}_{T}\left(\underline{\boldsymbol{e}}_{T},\underline{\boldsymbol{e}}_{T}\right)^{\frac{r+2-\bar{r}}{r}}$$

Stability and polynomial consistency are incompatible for k = 0!

Discrete stability hinges on the following result:

Theorem (Discrete Korn inequality)

Assume $k \ge 1$. Then, for all $\underline{v}_h \in \underline{U}_{h,0}^k$, letting $v_h \in \mathbb{P}^k(\mathcal{T}_h)^d$ be s.t. $(v_h)_{|T} \coloneqq v_T$ for all $T \in \mathcal{T}_h$,

$$\|\boldsymbol{v}_h\|_{L^r(\Omega)^d}+|\boldsymbol{v}_h|_{W^{1,r}(\mathcal{T}_h)^d}\lesssim \|\underline{\boldsymbol{v}}_h\|_{\varepsilon,r,h},$$

with $|\cdot|_{W^{1,r}(\mathcal{T}_h)^d}$ broken $W^{1,r}$ -seminorm.

Pressure-velocity coupling

The pressure-velocity coupling bilinear form $\mathbf{b}_h : \underline{U}_h^k \times \mathbb{P}^k(\mathcal{T}_h)$ is s.t.

$$\mathbf{b}_h(\underline{\boldsymbol{v}}_h,q_h)\coloneqq -\sum_{T\in\mathcal{T}_h}\int_T D_T^k\underline{\boldsymbol{v}}_T \ q_T$$

Lemma (Inf-sup stability)

Define the pressure space

$$P_h^k \coloneqq \left\{ q_h \in \mathbb{P}^k(\mathcal{T}_h) : \int_{\Omega} q_h = 0 \right\}.$$

Then it holds, for all $q_h \in P_h^k$,

$$\|q_h\|_{L^{r'}(\Omega)}\lesssim \sup_{\underline{\nu}_h\in\underline{U}_{h,0}^k,\,\|\underline{\nu}_h\|_{\varepsilon,r,h}=1}\mathrm{b}_h(\underline{\nu}_h,q_h).$$

Discrete problem

The discrete problem reads: Find $(\underline{u}_h, p_h) \in \underline{U}_{h,0}^k \times P_h^k$ s.t.

$$\begin{aligned} \mathbf{a}_{h}(\underline{\boldsymbol{u}}_{h},\underline{\boldsymbol{v}}_{h}) + \mathbf{b}_{h}(\underline{\boldsymbol{v}}_{h},p_{h}) &= \int_{\Omega} \boldsymbol{f} \cdot \boldsymbol{v}_{h} \quad \forall \underline{\boldsymbol{v}}_{h} \in \underline{\boldsymbol{U}}_{h,0}^{k}, \\ -\mathbf{b}_{h}(\underline{\boldsymbol{u}}_{h},q_{h}) &= 0 \qquad \forall q_{h} \in P_{h}^{k} \end{aligned}$$

Theorem (Well-posedness)

There exists a unique discrete solution $(\underline{u}_h, p_h) \in \underline{U}_{h,0}^k \times P_h^k$, and the following a priori bounds hold:

$$\begin{split} \|\underline{u}_{h}\|_{\varepsilon,r,h} &\lesssim \|f\|_{L^{r'}(\Omega)^{d}}^{\frac{1}{r-1}} + \|f\|_{L^{r'}(\Omega)^{d}}^{\frac{1}{r+1-\bar{r}}}, \\ \|p_{h}\|_{L^{r'}(\Omega)} &\lesssim \|f\|_{L^{r'}(\Omega)^{d}} + \|f\|_{L^{r'}(\Omega)^{d}}^{\frac{\bar{r}}{r+1-\bar{r}}}, \end{split}$$

with hidden multiplicative constants possibly depending on Ω , d, k, and the mesh regularity parameter.

Error estimate

Theorem (Error estimate)

Assume the regularity

$$\begin{split} \boldsymbol{u} &\in W^{1,r}(\Omega)^d \cap W^{k+2,r}(\mathcal{T}_{h})^d, \quad p \in W^{1,r'}(\Omega) \cap W^{(k+1)(\tilde{r}-1)}(\mathcal{T}_{h}), \\ \boldsymbol{\sigma}(\boldsymbol{\nabla}_{\mathbf{s}}\boldsymbol{u}) &\in W^{1,r'}(\Omega; \mathbb{R}^{d \times d}_{\mathrm{sym}}) \cap W^{(k+1)(\tilde{r}-1),r'}(\mathcal{T}_{h}; \mathbb{R}^{d \times d}_{\mathrm{sym}}). \end{split}$$

Then,

$$\begin{split} \|\underline{\boldsymbol{u}}_{h} - \underline{\boldsymbol{I}}_{h}^{k} \boldsymbol{u}\|_{\varepsilon, r, h} &\leq A h^{\frac{(k+1)(\tilde{r}-1)}{r+1-\tilde{r}}}, \\ \|p_{h} - \pi_{h}^{k} p\|_{L^{r'}(\Omega)} &\leq B h^{(k+1)(\tilde{r}-1)} + C h^{\frac{(k+1)(\tilde{r}-1)^{2}}{r+1-\tilde{r}}} \end{split}$$

with A, B, and C possibly depending on Ω , d, k, the mesh regularity parameter, and on bounded norms of u, p, and f.

Remark (Orders of convergence)

The order for the velocity is the same as for Leray-Lions problems. The asymptotic order for the pressure is $h^{(k+1)(r-1)^2}$ if r < 2, $\frac{k+1}{r-1}$ otherwise.

Numerical examples I

Convergence

- We assess the orders of convergence using a manufactured solution
- We take k = 1 and let r vary in $\{1.5, 1.75, \dots, 2.75\}$
- The regularity assumptions are mostly verified (except for r = 1.5, for which $\sigma(\nabla_{s} u) \notin W^{1,r'}(\Omega, \mathbb{R}^{d \times d}_{sym})$)
- We consider three families of meshes







Cartesian

Distorted triangular

Distorted quadrangular

Numerical examples II

Convergence



Figure: Convergence for shear-thinning fluids. The slopes indicate the expected order of convergence, i.e., $O_{vel}^1 = 2(r-1)$ and $O_{pre}^1 = 2(r-1)^2$ for $r \in \{1.5, 1.75, 2\}$.

Numerical examples III

Convergence



Figure: Convergence for shear-thickening fluids. The slopes indicate the expected order of convergence, i.e. $O_{vel}^1 = O_{pre}^1 = \frac{2}{r-1}$ for $r \in \{2.25, 2.5, 2.75\}$.

Lid-driven cavity I



Lid-driven cavity II



Figure: r = 1.25 (shear-thinning fluid)

Lid-driven cavity III



Figure: r = 2 (Newtonian fluid)

Lid-driven cavity IV



Figure: r = 2.75 (shear-thickening fluid)

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