# Convergence analysis of a numerical scheme for a tumour growth model



An Indian-Australian research partnership





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Numerical solutions of free boundary problems

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Numerical solutions of free boundary problems

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2 Discretisation

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cross-section of tumour spheroid

### Assumptions

- Cells and fluid exchange matter via the processes, cell division and cell death.
- Mass and momentum are conserved internally.
- No blood vessels. Limiting nutrient Oxygen, follows diffusion.

- Domain  $-0 < t < T, x \in \check{\Omega}(t) = (0, \check{\ell}(t)).$
- $\check{\ell}(t)$  tumour length, x = 0 tumour centre.
- $\check{\alpha}$  volume fraction of tumour cells,  $\check{u}$  cell velocity,  $\check{c}$  oxygen tension.

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•  $1 + (1/s_1)$ ,  $s_2 -$ maximal birth and death rates,  $s_3/s_4 -$ minimal death rate.

• Set 
$$f(\check{\alpha},\check{c}) = \frac{(1+s_1)(1-\check{\alpha})\check{c}}{1+s_1\check{c}} - \frac{s_2+s_3\check{c}}{1+s_4\check{c}}$$

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#### cell velocity (elliptic)

$$\begin{split} \frac{k\check{\alpha}\check{u}}{1-\check{\alpha}} &-\mu\frac{\partial}{\partial x}\left(\alpha\frac{\partial\check{u}}{\partial x}\right) = -\frac{\partial}{\partial x}\left(\check{\alpha}\mathcal{H}(\check{\alpha})\right),\\ \check{u}(t,0) &= 0, \quad \mu\frac{\partial\check{u}}{\partial x}(t,\ell(t)) = \mathcal{H}\left(\check{\alpha}(t,\ell(t))\right). \end{split}$$

- $\mu$  coefficient of viscosity of cell phase. k interfacial drag coefficient.
- Set  $\mathscr{H}(\check{\alpha}) = (\check{\alpha} \alpha^*)^+ / (1 \check{\alpha})^2, \ a^+ = \max(a, 0).$

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• Q – Maximum oxygen consumption rate.

- Domain  $-0 < t < T, x \in \check{\Omega}(t) = (0, \check{\ell}(t)).$
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#### boundary evolution

 $\check{\ell}'(t) = \check{u}(t,\check{\ell}(t)),$  $\check{\ell}(0) = 1.$ 

### Idea of extended model



- $\check{\ell}$  as the interface between  $\alpha > 0$  and  $\alpha = 0$ .
- velocity and oxygen tension extended by 0 and 1, respectively.

# Idea of threshold model



- $\ell$  as the interface between  $\alpha > 0$  and  $\alpha <= \alpha_{\text{thr}}$ .
- velocity and oxygen tension extended by 0 and 1, resp.
- α<sub>thr</sub> facilitates estimates on cell velocity and is required numerically.

A threshold solution (with threshold  $\alpha_{\text{thr}} \in (0,1)$ ) and domain  $D_T^{\text{thr}}$  of the threshold model in  $\mathscr{D}_T$  is a 4-tuple  $(\alpha, u, c, \Omega)$  such that:

- $0 < m_{11} \le \alpha_{|\Omega(t)} \le m_{12} < 1$  for all  $t \in [0, T]$ ,
- $m_{11} \le m_{01}, \, m_{12} \ge m_{02}$
- $c \ge 0$ ,

and the following hold:

### cell volume fraction

The volume fraction  $\alpha \in L^{\infty}(\mathscr{D}_T)$  is such that  $\forall \varphi \in \mathscr{C}^{\infty}_c([0,T) \times (0,\ell_m)),$ 

$$\int_{\mathscr{D}_T} (\alpha, u\alpha) \cdot \nabla_{t,x} \varphi \, \mathrm{d}t \, \mathrm{d}x + \int_{\Omega(0)} \varphi(0, x) \, \alpha_0 \, \mathrm{d}x + \int_{\mathscr{D}_T} (\alpha - \alpha_{\mathrm{thr}})^+ f(\alpha, c) \, \mathrm{d}x = 0$$

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### tumour boundary

The set  $D_T^{\text{thr}}$  is of the form

 $D_T^{\text{thr}} = \cup_{0 < t < T}(\{t\} \times \Omega(t)),$ 

where  $\Omega(t) = (0, \ell(t))$ , and we have  $\alpha \leq \alpha_{\text{thr}}$  on  $\mathscr{D}_T \setminus D_T^{\text{thr}}$ .

# Threshold solution

### cell velocity

• 
$$H^{1,u}_{\partial x}(D^{\text{thr}}_T) := \{ v \in L^2(D^{\text{thr}}_T) : \partial_x v \in L^2(D^{\text{thr}}_T)$$
  
and  $v(t,0) = 0 \ \forall t \in (0,T) \}.$   
•  $u \in H^{1,u}_{\partial x}(D^{\text{thr}}_T)$  and  $\forall v \in H^{1,u}_{\partial x}(D^{\text{thr}}_T)$ , satisfies  
 $\int_0^T a^t(u(t,\cdot), v(t,\cdot)) \, \mathrm{d}t = \int_0^T \mathcal{L}^t(v(t,\cdot)) \, \mathrm{d}t,$  (1)

where  $a^t : H^1(\Omega(t)) \times H^1(\Omega(t)) \to \mathbb{R}$  is a bilinear form and  $\mathcal{L}^t : H^1(\Omega(t)) \to \mathbb{R}$  is a linear form as follows:

$$a^{t}(u,v) = k \left(\frac{\alpha}{1-\alpha}u,v\right)_{\Omega(t)} + \mu \left(\alpha \partial_{x}u,\partial_{x}v\right)_{\Omega(t)} \text{ and }$$
(2)

$$\mathcal{L}^{t}(v) = (\mathscr{H}(\alpha), \partial_{x}v)_{\Omega(t)}.$$
(3)

Extend u to  $\mathscr{D}_T$  by setting  $u|_{\mathscr{D}_T \setminus \overline{D}_T^{\text{thr}}} := 0.$ 

## Threshold solution

#### oxygen tension

• 
$$\begin{aligned} H^{1,c}_{\partial x}(D^{\text{thr}}_T) &:= \{ v \in L^2(D^{\text{thr}}_T) : \partial_x v \in L^2(D^{\text{thr}}_T) \\ \text{and } v(t,\ell(t)) &= 0 \ \forall t \in (0,T) \}. \end{aligned} \\ \bullet \ c-1 \in H^{1,c}_{\partial x}(D^{\text{thr}}_T) \text{ satisfies,} \\ - \int_{D^{\text{thr}}_T} c \partial_t v \, dx \, dt + \lambda \int_{D^{\text{thr}}_T} \partial_x c \partial_x v \, dx \, dt - \int_{\Omega(0)} c_0(x) v(0,x) \, dx \\ - Q \int_{D^{\text{thr}}_T} \alpha c v \, dx \, dt = 0, \end{aligned}$$

 $\begin{array}{l} \forall v \in H^{1,c}_{\partial x}(D^{\mathrm{thr}}_T) \text{ such that } \partial_t v \in L^2(D^{\mathrm{thr}}_T). \text{ Extend } c \text{ to } \mathscr{D}_T \text{ by setting } c \big|_{\mathscr{D}_T} \backslash \overline{D^{\mathrm{thr}}_T} := 1. \end{array}$ 

Numerical solutions of free boundary problems

- To obtain a lower bound strictly greater than zero for  $\alpha$ .
- Facilitates bounded variation estimates on  $\alpha$ .
- To obtain supremum norm bounds on u and  $\partial_x u$ .

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### An unavoidable disadvantage

- Residual volume fraction creates spurious growth outside the tumour domain.
- Essential from numerical vantage point.
- Modified source term  $(\alpha \alpha_{thr})^+ f(\alpha, c)$  eliminates spurious growth.
- As  $\alpha_{\text{thr}} \to 0$ ,  $(\alpha \alpha_{\text{thr}})^+ f(\alpha, c)$  approaches  $\alpha f(\alpha, c)$ .



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### Discrete scheme

- Space:  $0 = x_0 < \dots < x_J = \ell_m$ , Time:  $0 = t_0 < \dots t_n = T$
- Uniform discretisation:  $\delta = t_{n+1} t_n$ ,  $h = x_{j+1} x_j$ .

#### $\operatorname{scheme}$

- volume fraction:  $\alpha_h^n$  upwind finite volume scheme.
- Set  $\ell_h^n = \min\{x_j : \alpha_j^n < \alpha_{\text{thr}} \text{ on } (x_j, \ell_m)\}$  and  $\Omega_h^n := (0, \ell_h^n).$
- Conforming Lagrange  $\mathbb{P}^1$ -FEM to obtain  $u_{h|\Omega_h^n}^n$ , and set  $u_h^n = 0$  outside  $\Omega_h^n$ .
- Time-implicit mass lumped  $\mathbb{P}^1$ -FEM to obtain  $c_{h|\Omega_h^n}^n$ , and set  $c_h^n = 1$  outside  $\Omega_h^n$ .

### Definition (Time-reconstruct)

For a family of functions  $(f_h^n)_{\{0 \le n < N\}}$  on a set X, define the time-reconstruct  $f_{h,\delta}: (0,T) \times X \to \mathbb{R}$  as  $f_{h,\delta}:=f_h^n$  on  $[t_n, t_{n+1})$  for  $0 \le n < N$ .

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#### Definition (Discrete solution)

The 4-tuple  $(\alpha_{h,\delta}, u_{h,\delta}, c_{h,\delta}, \ell_{h,\delta})$ , where  $\alpha_{h,\delta}, u_{h,\delta}, c_{h,\delta}$ , and  $\ell_{h,\delta}$  are the respective time-reconstructs corresponding to the families  $(\alpha_h^n)_n, (u_h^n)_n, (c_h^n)_n$ , and  $(\ell_h^n)_n$  is called the discrete solution.

# Why a mixed numerical scheme?

### Finite volume method

- Respects mass conservation property at the discrete level.
- Upwind flux (+ CFL) yields a stable scheme.
- FVM significant numerical diffusion.
- Large error in locating  $\ell_h^n$  as the boundary where  $\alpha_h^n$  becomes 0.
- Solution: Locate  $\ell_h^n$  as  $\min\{x_j : \alpha_h^n < \alpha_{\text{thr}} \text{ on } (x_j, \ell_m]\}.$



### Finite element methods

- Velocity equation elliptic, oxygen tension equation parabolic.
- Unknown Lagrange  $\mathbb{P}^1$ -FEM boundary nodes of  $(x_j, x_{j+1})$ , and easy to compute the upwind flux.
- Mass lumping in oxygen tension equation is crucial to obtain  $L^\infty$  bounds.
- Time-implicitness yields stability in  $L^2(0,T; H^1(0,\ell_m))$ .





### 3 Main Theorem

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# $\hat{u}_{h,\delta}$ : continuous modification of $u_{h,\delta}$



Figure: The left-hand side plot illustrates the discontinuous function  $u_{h,\delta}$ and the right-hand side plot illustrates the continuous modification  $\hat{u}_{h,\delta}$ .

### Few notations

#### Mass lumping operator

- Set  $\tilde{\chi}_j = (x_j h/2, x_j + h/2)$ ,  $S_{h,ML}$  piecewise constant functions on  $\tilde{\chi}_j$ .
- Mass lumping operator:  $\Pi_h : \mathscr{C}^0([0,L]) \to \mathcal{S}_{h,ML}$  such that  $\Pi_h w = \sum_{j=0}^J w(x_j) \mathbf{1}_{\widetilde{\chi}_j}$ .
- Set  $\Pi_{h,\delta}c_{h,\delta}$  by  $\Pi_{h,\delta}c_{h,\delta}(t,\cdot) := \Pi_h(c_{h,\delta}(t,\cdot)).$



Numerical solutions of free boundary problems

### Time-dependent spaces

$$\begin{split} L^2_c(0,T;H^1(0,\ell_m)) &:= \{f \in L^2(0,T;H^1(0,\ell_m)) \,:\, f(t,\ell(t)) = 0 \\ & \text{for a.e. } t \in [0,T] \}, \\ L^2_u(0,T;H^1(0,\ell_m)) &:= \{f \in L^2(0,T;H^1(0,\ell_m)) \,:\, f(t,0) = 0 \\ & \text{for a.e. } t \in [0,T] \}. \end{split}$$

Numerical solutions of free boundary problems

#### Theorem (compactness)

Let the properties stated below be true.

- The initial volume fraction  $\alpha_0$  belongs to  $BV(0, \ell_m)$  and  $0 < m_{01} \le \alpha_0 \le m_{02} < 1$ , where  $m_{01}$  and  $m_{02}$  are constants.
- The discretisation parameters h and  $\delta$  satisfy the following conditions:

$$\rho \mathscr{C}_{CFL} \leq \frac{\delta}{h} \leq \mathscr{C}_{CFL} := \frac{\sqrt{a_*}\mu}{2\ell_m} \frac{|1-a^*|^2}{|a^*-\alpha^*|} \text{ and}$$
$$\delta < \min\left(\frac{1-\rho}{s_2}, \frac{2(1-\rho)}{1+s_2}\right),$$

where  $\rho$ ,  $a_*$  and  $a^*$  are constants chosen such that  $\rho < 1$ ,  $0 < a_* < m_{01}$ , and  $0 < m_{02} < a^*$ .

#### Theorem (compactness)

Then, there exists a finite time  $T_* = T_*(\rho, a_*, a^*)$ , a subsequence of the family of functions  $\{(\alpha_{h,\delta}, \hat{u}_{h,\delta}, c_{h,\delta}, \ell_{h,\delta})\}_{h,\delta}$  and a 4-tuple of functions  $(\alpha, \hat{u}, c, \ell)$  such that

- $\alpha \in BV(\mathscr{D}_{T_*})$
- $c \in L^2_c(0, T_*; H^1(0, \ell_m))$
- $\hat{u} \in L^2_u(0, T_*; H^1(0, \ell_m))$
- $\ell \in BV(0,T_*)$

with  $\mathscr{D}_{T_*} = (0, T_*) \times (0, \ell_m)$  and as  $h, \delta \to 0$ ,

- $\alpha_{h,\delta} \to \alpha$  almost everywhere and in  $L^{\infty}$  weak<sup>\*</sup> on  $\mathscr{D}_{T_*}$ ,
- $\Pi_{h,\delta}c_{h,\delta} \to c$  strongly in  $L^2(\mathscr{D}_{T_*})$  and  $\partial_x c_{h,\delta} \rightharpoonup \partial_x c$  weakly in  $L^2(\mathscr{D}_{T_*})$ ,
- $\widehat{u}_{h,\delta} \rightharpoonup \widehat{u}$  and  $\partial_x \widehat{u}_{h,\delta} \rightharpoonup \partial_x \widehat{u}$  weakly in  $L^2(\mathscr{D}_{T_*})$ , and
- $\ell_{h,\delta} \to \ell$  almost everywhere in  $(0,T_*)$ .

#### Theorem (convergence)

Let  $(\alpha, \hat{u}, c, \ell)$  be the limit provided by the compactness theorem. Define  $\Omega(t) := (0, \ell(t))$  and the threshold domain

 $D_{T_*}^{\mathrm{thr}} := \{(t,x): x < \ell(t), t \in (0,T_*)\}$ 

and let  $u := \hat{u}$  on  $D_{T_*}^{\text{thr}}$  and u := 0 on  $\mathscr{D}_{T_*} \setminus D_{T_*}^{\text{thr}}$ . Then,  $(\alpha, u, c, \Omega)$  is a threshold solution with  $T = T_*$ .

# Idea for proof of Main Theorem I

- Proof is through inductive arguments on time-step, n.
- Fix two constants  $a^* \in (\max(\alpha^*, m_{02}), 1)$  and  $a_* \in (0, \min(\alpha_{\text{thr}}, m_{01})).$
- The time of existence  $T_*$  on  $a^*$  and  $a_*$ , and is explicitly provided by Theorem (well-posedness).

#### Theorem (well-posedness)

For all  $n \in \mathbb{N}$  such that  $t_n \leq T_*$ ,  $\alpha_{h,\delta}(t_n, \cdot)$ ,  $u_{h,\delta}(t_n, \cdot)$ , and  $c_{h,\delta}(t_n, \cdot)$  are well defined, and it holds:

• 
$$a_* < \alpha_{h,\delta}(t_n,\cdot)|_{\Omega^n_h} < a^*,$$

• 
$$0 \le c_{h,\delta}(t_n, \cdot)_{|(0,\ell_m)} \le 1.$$

• Necessary compactness results proved using supremum norm bounds from the Theorem (well-posedness).

### Proof of well-posedness Theorem

#### Step 1: Energy estimate of $\tilde{u}_h^n$

There exists a unique solution  $\widetilde{u}_h^n$  to discrete weak form of velocity equation in  $\Omega_h^n$  and it satisfies the following estimates:

$$\left\| \sqrt{\alpha_{h,\delta}(t_n,\cdot)} \partial_x \widetilde{u}_h^n \right\|_{0,\Omega_h^n} \leq \frac{\sqrt{\ell_m} |a^* - \alpha^*|}{\mu |1 - a^*|^2} \text{ and} \\ \left\| \frac{\sqrt{\alpha_{h,\delta}(t_n,\cdot)} \widetilde{u}_h^n}{\sqrt{1 - \alpha_{h,\delta}(t_n,\cdot)}} \right\|_{0,\Omega_h^n} \leq \sqrt{\frac{\ell_m}{k\mu}} \frac{|a^* - \alpha^*|}{|1 - a^*|^2}.$$

- Keeping the coefficients yields optimal estimates, which improves the existence time  $T_*$ .
- Estimate on  $\partial_x \widetilde{u}_h^n$  yields

$$|u_{h,\delta}(t_n,\cdot)||_{L^{\infty}(0,\ell_m)} \le \frac{\ell_m}{\sqrt{a_*\mu}} \frac{|a^* - \alpha^*|}{|1 - a^*|^2}.$$
(3)

### Proof of well-posedness Theorem

### Step 3: BV and $L^{\infty}$ bound on $\partial_x u_{h,\delta}(t_n,\cdot)$

### It holds:

$$\begin{split} ||\mu\alpha_{h,\delta}(t_{n},\cdot)\partial_{x}u_{h,\delta}(t_{n},\cdot) - \mathscr{H}(\alpha_{h,\delta}(t_{n},\cdot))||_{BV(0,\ell_{m})} \\ &\leq \ell_{m}\sqrt{\frac{k}{\mu}}\frac{|a^{*}-\alpha^{*}|}{|1-a^{*}|^{5/2}}, \\ ||(\mu\alpha_{h,\delta}(t_{n},\cdot)\partial_{x}u_{h,\delta}(t_{n},\cdot))^{-}||_{L^{\infty}(0,\ell_{m})} \leq \ell_{m}\sqrt{\frac{k}{\mu}}\frac{|a^{*}-\alpha^{*}|}{|1-a^{*}|^{5/2}}, \text{ and} \\ ||\mu\alpha_{h,\delta}(t_{n},\cdot)\partial_{x}u_{h,\delta}(t_{n},\cdot)||_{L^{\infty}(0,\ell_{m})} \leq \ell_{m}\sqrt{\frac{k}{\mu}}\frac{|a^{*}-\alpha^{*}|}{|1-a^{*}|^{5/2}} \\ &+ \frac{a^{*}(a^{*}-\alpha^{*})}{(1-a^{*})^{2}}. \end{split}$$

Step 3:  $L^{\infty}$  bound on  $\alpha_{h,\delta}(t_n,\cdot)$ 

There exists  $T_* > 0$  such that if  $n+1 \leq N_* := T_*/\delta$ , then

$$a_* \le \min_{j: x_j \in \Omega_h^{n+1}} \alpha_j^{n+1} \le \max_{0 \le j \le J-1} \alpha_j^{n+1} \le a^*.$$

### Step 4: $\overline{L^{\infty}}$ bound on $c_{\overline{h,\delta}}(t_n,\cdot)$

The discrete weak form corresponding to the oxygen tension equation has a unique solution  $\tilde{c}_h^{n+1}$  in  $\Omega_h^n$ , and it holds  $0 \leq \tilde{c}_h^{n+1} \leq 1$ .



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### Check-list of compactness estimates

- a uniform  $L^2(0, T_*; H^1(0, \ell_m))$  estimate for the family  $\{c_{h,\delta}\}_{h,\delta}$  weak  $L^2(0, T_*; H^1(0, \ell_m))$  convergence.
- **2** a uniform spatial and temporal BV estimate for the family  $\{\alpha_{h,\delta}\}_{h,\delta}$  strong  $L^p(\mathscr{D}_{T_*})$  convergence.
- (a uniform BV estimate for the family  $\{\ell_{h,\delta}\}_{h,\delta}$  strong  $L^p(0,T_*)$  convergence.
- the family  $\{\Pi_{h,\delta}c_{h,\delta}\}_{h,\delta}$  is relatively compact in  $L^2(\mathscr{D}_{T_*})$  strong  $L^2(\mathscr{D}_{T_*})$  convergence.
- a uniform  $L^2(0,T_*;H^1(0,\ell_m))$  estimate for the family  $\{\widehat{u}_{h,\delta}\}_{h,\delta}$  weak  $L^2(0,T_*;H^1(0,\ell_m))$  convergence.

# Relative compactness of $\{\Pi_h c_{h,\delta}\}_{h,\delta}$

### Definition

Auxiliary function Define  $\hat{c}_{h,\delta} := c_{h,\delta} - 1$ . For a fixed  $\epsilon > 0$ , define the auxiliary function  $\varphi_{h,\epsilon}^n : [0, \ell_m] \to [0, 1]$  by

$$\varphi_{h,\epsilon}^{n}(x) = \begin{cases} 1 & 0 \le x \le \ell_{h}^{n} - \epsilon, \\ (\ell_{h}^{n} - x)/\epsilon & \ell_{h}^{n} - \epsilon < x \le \ell_{h}^{n}, \\ 0 & \ell_{h}^{n} < x \le \ell_{m}. \end{cases}$$



• The mass lumped function can be split into

 $\Pi_{h,\delta}\widehat{c}_{h,\delta} = \Pi_{h,\delta}(\widehat{c}_{h,\delta}\varphi_{h,\epsilon}) + \Pi_{h,\delta}(\widehat{c}_{h,\delta}(1-\varphi_{h,\epsilon})),$ 

where  $\varphi_{h,\epsilon} = \varphi_{h,\epsilon}^n$  on  $[t_n, t_{n+1})$  for  $0 \le n \le N_* - 1$ .

• The second term can be bounded by:

$$||\Pi_{h,\delta}(\widehat{c}_{h,\delta}(1-\varphi_{h,\epsilon}))||_{L^2(\mathscr{D}_{T_*})} \leq \sqrt{T_*\epsilon}.$$

#### Theorem

The family of functions  $\{\Pi_{h,\delta}(\varphi_{h,\epsilon}\widehat{c}_{h,\delta})\}_{h,\delta}$  is relatively compact in  $L^2(\mathscr{D}_{T_*})$ .

• Proof follows from the Discrete Aubin - Simon Theorem.

#### Theorem

The family of functions  $\{\Pi_{h,\delta}c_{h,\delta}\}_{h,\delta}$  is relatively compact in  $L^2(\mathscr{D}_{T_*})$ .

• Proof follows from the fact  $\epsilon > 0$ ,

 $\{\Pi_{h,\delta}\widehat{c}_{h,\delta}\}_{h,\delta} \subset \{\Pi_{h,\delta}(\varphi_{h,\epsilon}\widehat{c}_{h,\delta})\}_{h,\delta} + B_{L^2(\mathscr{D}_{T_*})}\left(0;\sqrt{T_*\epsilon}\right).$ 



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The domains

$$A_{h,\delta} := \{ (t,x) : x < \ell_{h,\delta}(t), t \in (0,T_*) \}$$

converge to

$$D_{T_*}^{\text{thr}} := \{(t, x) : x < \ell(t), t \in (0, T_*)\}.$$

- **②** The limit function  $\alpha$  satisfies the weak form (volume fraction) with  $T = T_*$ .
- **③** The restricted limit function  $\hat{u}_{|D_{T_*}^{\text{thr}}}$  satisfies weak form (cell velocity) with  $T = T_*$ .
- The limit function  $c_{|D_{T_*}^{\text{thr}}}$  satisfies weak form (oxygen tension) with  $T = T_*$ .



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### Parameters

- $k = 1, \mu = 1, Q = 0.5, s_1 = 10 = s_4, s_2 = 0.5 = s_3, \alpha^* = 0.8^1.$
- The bounds of the cell volume fraction are set to be  $a_* = 0.4$  and  $a^* = 0.82$ .
- The extended domain length  $\ell_m$  is set as 10.
- The threshold value is taken as  $\alpha_{thr} = 0.1$ .

• 
$$\rho = 0.1, \ \delta = 1E - 3 \text{ and } h = 5E - 2.$$

- Set  $T_* = 50$ .
- Predicted time by compactness theorem: 1E 7 to 1E 1.

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<sup>&</sup>lt;sup>1</sup>Breward, C.J.W., Byrne, H.M. and Lewis, C.E., 2002. The role of cell-cell interactions in a two-phase model for avascular tumour growth. *J. of Math. Bio.*, 45(2), pp. 125-152.



- Sufficiency of compactness and convergence theorems Existence of solutions beyond  $T_*$  is possible.
- Convergence theorem guarantees existences of a domain  $D_T^{\alpha_{\text{thr}}}$ . However, it is not known whether it is unique.
- Framework can be extended to similar problems and models.
- Higher dimensional study (on going work).

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# Thanks

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