Design and analysis of an extended virtual element method

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methods for numerical

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Principles of mesh-based schemes

Model problem

$$
\begin{cases}\n-\Delta u = f & \text{in } \Omega, \\
u = 0 & \text{on } \partial \Omega.\n\end{cases}
$$

Mesh-based methods: cut domain in small pieces (elements/cells), approximate the solution with piecewise polynomial functions on these pieces.

Typical error estimate: for a method of order $k \geq 1$,

$$
||u - u_h||_{H^1} \le Ch^k ||u||_{H^{k+1}}
$$

◦ Comes from the local approximation properties of polynomial functions.

 \circ Requires smoothness of u .

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Failure of smoothness: even in non-challenging situation (any non-convex domain!).

 $u\not\in H^2$, so poor polynomial approximation...

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Failure of smoothness: even in non-challenging situation (any non-convex domain!).

$$
u = r^{\frac{2}{3}} \sin\left(\frac{2}{3}(\theta - \frac{\pi}{2})\right) + u_r \text{ with } u_r \in H^2.
$$

Extend the local polynomial spaces with the singular part of the solution.

Caveat: far from the singularity, the "singular" part is actually smooth, and too close to polynomials...

Refs.:

◦ FEM: [\[Melenk and Babuˇska, 1996,](#page-31-0) [Babuˇska and Melenk, 1997,](#page-30-0) [Belytschko and Black, 1999,](#page-30-1) [Laborde et al., 2005,](#page-31-1) [Chin et al., 2017\]](#page-31-2) etc.

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Polytopal methods and virtual elements? I

- Finite Element methods (FEM) have been the golden standard for decades, but their mesh lack flexibility with: local refinement, representation of complex geometries, etc.
- Polytopal methods are inherited from FEM but are applicable on meshes made of generic polygons/polyhedra.

Refs.: [Beirão da Veiga et al., 2014, Beirão da Veiga et al., 2017, [Di Pietro and Droniou, 2020\]](#page-31-3) etc.

- Virtual Element Method is a polytopal method based on approximation spaces such that:
	- □ Virtual functions are not fully known,
	- \Box Unisolvent degrees of freedom (DOFs) can be identified for the spaces,
	- \square Certain projections of virtual functions/gradients can be explicitly computed from the DOFs.

Local standard space: Fix $k \ge 1$, $l = \max(0, k - 2)$ and define, for each element E,

$$
V_{k,h}(E) := \{ v_h \in H^1(E) : \Delta v_h \in \mathbb{P}_l(E),
$$

$$
v_{h|\partial E} \in C^0(\partial E), v_{h|\varepsilon} \in \mathbb{P}_k(e) \ \forall e \subset \partial E \}.
$$

Contains polynomials: $\mathbb{P}_k(E) \subset V_{k,h}(E)$.

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$$

Contains polynomials: $\mathbb{P}_k(E) \subset V_{k,h}(E)$.

Local extended space: with $\Psi \subset H^1(\Omega)$ the space of singularities of the solution,

$$
V_{k,h}^{\Psi}(E) := \{ v_h \in H^1(E) : \Delta v_h \in \mathbb{P}_l(E) + \Delta \Psi_{|E},
$$

$$
v_{h|\partial E} \in C^0(\partial E), v_{h|e} \in \mathbb{P}_k(e) + \Psi_{|e} \ \forall e \subset \partial E \}.
$$

Contains polynomials and singular function: $\mathbb{P}_k(E) + \Psi_{|E} \subset V_{k,h}^{\Psi}(E)$.

$$
V_{k,h}^{\Psi}(E) := \{ v_h \in H^1(E) : \Delta v_h \in \mathbb{P}_l(E) + \Delta \Psi_{|E},
$$

$$
v_{h|\partial E} \in C^0(\partial E), v_{h|e} \in \mathbb{P}_k(e) + \Psi_{|e} \ \forall e \subset \partial E \}.
$$

Decomposition of boundary space: write $\mathbb{P}_k(e) + \Psi_{|e} = \mathbb{P}_k(e) \oplus \mathfrak{P}_e$ and set

$$
\mathbb{P}_{k-2}^{\mathfrak{P}}(e) = \mathbb{P}_{k-2}(e) \oplus \mathfrak{P}_e.
$$

Degrees of freedom: for $v_h \in V_{k,h}^{\Psi}(E)$, (D1) $v_h(x_V)$ for each vertex $V \in \partial E$; $\pmb{(D2)}\ \Pi_{k-2,e}^{\mathfrak{P}} v_h$, projection on $\mathbb{P}_{k-2}^{\mathfrak{P}}(e)$, for each edge $e\subset\partial E;$ $\left(\textbf{D3}\right) \Pi_{l,E}^{\Delta} v_h$, projection on $\mathbb{P}_{l}(E) + \Delta \Psi_{|E}$.

 $\mathbb{P}_{k-2}(e) + \Psi_{|e}$ would be natural, but does not ensure unisolvence...

Extended elliptic projector: $\Pi_{k,E}^{\nabla,\Psi}: V_{k,h}^{\Psi}(E) \to \mathbb{P}_k(E) + \Psi_{|E}$, computable from the DOFs, defined by

$$
\int_{E} \nabla(\Pi_{k,E}^{\nabla,\Psi} v_h) \cdot \nabla q = \int_{E} \nabla v_h \cdot \nabla q \quad q \in \mathbb{P}_k(E) + \Psi_{|E},
$$

$$
\int_{E} \Pi_{k,E}^{\nabla,\Psi} v_h = \int_{E} v_h.
$$

Remark: $\Pi_{k,E}^{\nabla,\Psi}v = v$ for all $v \in \mathbb{P}_k(E) + \Psi_{|E}$.

Global space: $V_{k,h,0}^{\Psi}$ obtained by gluing the local spaces, and

$$
V_{k,h,0}^{\Psi} := \{ v_h \in H_0^1(\Omega) : v_{h|E} \in V_{k,h}^{\Psi}(E) \quad \forall E \in \Omega_h \}.
$$

Extended virtual element scheme II

Bilinear form: consistent part plus stabilisation.

$$
a_h(u_h, v_h) := \sum_{E \in \Omega_h} a_E(u_h, v_h)
$$

with

$$
a_E(u_h, v_h) = \int_E \nabla \Pi_{k,E}^{\nabla, \Psi} u_h \cdot \nabla \Pi_{k,E}^{\nabla, \Psi} v_h + S_E(u_h, v_h)
$$

and S_E such that

$$
S_E(u_h, u_h) = h_E^{-2} \|\Pi_{l,E}^{\Delta}(u_h - \Pi_{k,E}^{\nabla,\Psi} u_h)\|_{L^2(E)}^2 + h_E^{-1} \|u_h - \Pi_{k,E}^{\nabla,\Psi} u_h\|_{L^2(\partial E)}^2.
$$

Scheme: Find $u_h \in V_{k,h,0}^{\Psi}$ such that

$$
a_h(u_h, v_h) := \sum_{E \in \Omega_h} \int_E f \Pi_{l,E}^{\Delta} v_h \quad \forall v_h \in V_{k,h,0}^{\Psi}.
$$

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Theorem (Discrete Energy Error)

Let $u = u^r + \psi$ solution to PDE with $\psi \in \Psi$ and $u^r \in H^{k+1}(\Omega_h)$. Under standard mesh regularity assumption:

$$
||u_h - \mathcal{I}_{k,h}u||_{a,h} \lesssim h^k |u^r|_{H^{k+1}(\Omega_h)},
$$

where $\|\cdot\|_{a,h}$ is the energy norm associated with a_h and $\mathcal{I}_{k,h}u = \mathcal{I}_{k,h}u^r + \psi$ with $\widehat{\mathcal{I}}_{k,h}$ standard VEM interpolant of u^r .

Consistency analysis – The Virtual Element Approach

$$
\sum_{E \in \Omega_h} \int_E -\Delta u \Pi_{l,E}^{\Delta} v_h - a_h(\mathcal{I}_{k,h} u, v_h)
$$
\n
$$
= \int_{\Omega} -\Delta u v_h + \mathcal{O}(h^k) - \sum_{E \in \Omega_h} \int_E \nabla (\Pi_{k,E}^{\nabla, \Psi} \mathcal{I}_{k,h} u) \cdot \nabla \Pi_{k,E}^{\nabla, \Psi} v_h + \text{stab}
$$
\n
$$
= \int_{\Omega} \nabla (u - \Pi_{k,E}^{\nabla, \Psi} \mathcal{I}_{k,h} u) \cdot \nabla v_h + \mathcal{O}(h^k) + \text{stab}
$$
\n
$$
\leq \|\nabla (u - \Pi_{k,E}^{\nabla, \Psi} \mathcal{I}_{k,h} u) \|_{L^2} \|\nabla v_h\|_{L^2} + \mathcal{O}(h^k) + \text{stab}
$$

To conclude:

- \circ Approximation properties of $\mathcal{I}_{k,h}u$.
- Boudedness $\|\nabla v_h\|_{L^2} \leq C \|v_h\|_{a,h}$ (norm on DOFs of v_h).

Difficult for standard VEM, not known for extended VEM...

Refs: [\[Benvenuti et al., 2019,](#page-30-4) [Artioli and Mascotto, 2021\]](#page-30-5).

Circumvents these issues by adopting a fully discrete approach: do not introduce virtual functions, express everything in terms of computable quantities.

Refs: Enriched Hybrid High-Order: [\[Yemm, 2022,](#page-31-4) [Yemm, 2024\]](#page-31-5).

Consistency analysis – The Fully Discrete Approach II

Manipulate source term: introduce elliptic projector, perform IBP, use continuity of fluxes:

$$
- \sum_{E \in \Omega_h} \int_E \Delta u \Pi_{l,E}^{\Delta} v_h
$$

=
$$
- \sum_{E \in \Omega_h} \int_E \Delta u (\Pi_{l,E}^{\Delta} v_h - \Pi_{k,E}^{\nabla, \Psi} v_h) + \sum_{E \in \Omega_h} \int_E \nabla u \cdot \nabla \Pi_{k,E}^{\nabla, \Psi} v_h
$$

+
$$
\sum_{E \in \Omega_h} \langle \nabla u \cdot \mathbf{n}, v_h - \Pi_{k,E}^{\nabla, \Psi} v_h \rangle_{\partial E}.
$$

Property of projectors: for all $z \in \mathbb{P}_k(E) + \Psi_{|E}$, we have $\Delta z \in \mathbb{P}_{l}(E) + \Delta \Psi_{|E}$ so

$$
-\int_{E} \Delta z \left(\prod_{k} \mathcal{L}v_{h} - \Pi_{k,E}^{\nabla,\Psi}v_{h}\right) + \langle \nabla z \cdot \mathbf{n}, v_{h} - \Pi_{k,E}^{\nabla,\Psi}v_{h} \rangle_{\partial E}
$$

$$
\stackrel{\text{IBP}}{=} \int_{E} \nabla z \cdot \nabla \left(v_{h} - \Pi_{k,E}^{\nabla,\Psi}v_{h}\right) \stackrel{\text{def. } \Pi_{k,E}^{\nabla,\Psi}}{=} 0.
$$

Consistency analysis – The Fully Discrete Approach III

Eliminate singular part of u: recalling that $u = u^r + \psi$, take $z = q_E + \psi \in \mathbb{P}_k(E) + \Psi_{|E}$ and subtract:

$$
- \sum_{E \in \Omega_h} \int_E \Delta u \Pi_{l,E}^{\Delta} v_h = \sum_{E \in \Omega_h} \int_E \nabla u \cdot \nabla \Pi_{k,E}^{\nabla, \Psi} v_h
$$

$$
- \sum_{E \in \Omega_h} \int_E \Delta (u^r - q_E) (\Pi_{l,E}^{\Delta} v_h - \Pi_{k,E}^{\nabla, \Psi} v_h)
$$

$$
+ \sum_{E \in \Omega_h} \int_{\partial E} \nabla (u^r - q_E) \cdot \mathbf{n} (v_h - \Pi_{k,E}^{\nabla, \Psi} v_h).
$$

- $\circ\,$ Last two terms are $\mathcal{O}(h^k)$ by regularity of u^r and polynomial approximation.
- \circ First term combines with consistent term in $a_h(\mathcal{I}_{k,h}u, v_h)$:

$$
\int_{E} \nabla (\Pi_{k,E}^{\nabla,\Psi} \mathcal{I}_{k,h} u) \cdot \nabla \Pi_{k,E}^{\nabla,\Psi} v_h = \int_{E} \nabla \mathcal{I}_{k,h} u \cdot \nabla \Pi_{k,E}^{\nabla,\Psi} v_h
$$

to create $u - \mathcal{I}_{k,h} u = u^r - \widehat{\mathcal{I}}_{k,h} u^r$ (since $\psi = \mathcal{I}_{k,h} \psi$), which is $\mathcal{O}(h^k)$.

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L-shaped domain $\Omega = (-1, 1)^2 \setminus [0, 1)^2$

Solution: take $u = sin(\pi x) sin(\pi y) + \psi$ with singularity at re-entrant corner

$$
\psi(r,\theta) = r^{\frac{2}{3}} \sin(\frac{2}{3}(\theta - \frac{\pi}{2})).
$$

Mesh:

L-shaped domain $\Omega = (-1, 1)^2 \setminus [0, 1)^2$

Fractured domain $\Omega = (0, 1)^2 \setminus ([0, 1) \times \{0\})$

Solution: take $u = sin(\pi x) sin(\pi y) + \psi$ with singularity at re-entrant corner

$$
\psi(r,\theta) = r^{\frac{1}{2}}\sin(\frac{1}{2}\theta).
$$

Mesh:

Fractured domain $\Omega = (0, 1)^2 \setminus ([0, 1) \times \{0\})$

- Polytopal method, benefitting from the flexibility of general polygonal/polyhedral meshes.
- Extended: includes a singularity space in the design, to better reproduce singular solutions.
- Recovers optimal convergence for problems with re-entrant corners and cracks. But also improves problems with highly oscillatory solutions (PhD L. Yemm).
- First complete analysis, circumvents the issues of a virtual-element based analysis by using a fully discrete approach (only based on DOFs, not virtual functions).

N NEMESIS

New generation methods for numerical simulations

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Thank you for your attention!

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