

A nonconforming Virtual Element Method for a biharmonic problem on polygonal meshes

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Outline

- 1 Continuous problem
- 2 Non-conforming VEM spaces
- 3 Discrete VEM problem
- 4 Numerical results
- 5 Conclusions

Introduction

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FV: [Eymard, Gallouet, Herbin, Linke, 2012], HHO: [Chave, Di Pietro, Marche, Pigeonneau, 2016]

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Goal: Design a “weaker” **nonconforming** Virtual Element Method (VEM) for the discretization of the biharmonic problem on **polygonal meshes**.

The continuous problem

Let $\Omega \subset \mathbb{R}^2$ convex polygonal domain and $f \in L^2(\Omega)$

$$\begin{aligned} D\Delta^2 u &= f && \text{in } \Omega \\ u = \partial_n u &= 0 && \text{on } \Gamma = \partial\Omega \end{aligned}$$

$D = \frac{Et^3}{12(1-\nu^2)}$ bending rigidity, t thickness, E Young modulus, ν Poisson's ratio.

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Weak formulation: Find $u \in V = \{v \in H^2(\Omega) : v = \partial_n v = 0 \text{ on } \Gamma\}$ s.t.

$$a(u, v) = F(v) \quad \forall v \in V$$

where

$$a(u, v) = D \int_{\Omega} \nu \Delta u \Delta v + (1 - \nu)(u_{,ij} v_{,ij}) dx \quad F(v) = \int_{\Omega} f v dx.$$

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Setting $\|\cdot\|_V := |\cdot|_{2,\Omega}$ there exist $\alpha, M > 0$ such that

$$\begin{aligned} a(v, v) &\geq \alpha \|v\|_V^2 & \forall v \in V \\ |a(u, v)| &\leq M \|u\|_V \|v\|_V & \forall u, v \in V. \end{aligned}$$

Hence, there exists a unique solution $u \in V$.

Let $\sigma_{ij}(u) = \lambda(u_{,11} + u_{,22})\delta_{ij} + \mu u_{,ij}$ with Lamé parameters $\lambda = D\nu$, $\mu = D(1 - \nu)$. We set:

$$M_{nn}(u) = \sum_{i,j} \sigma_{ij} n_i n_j \quad (\text{normal bending moment})$$

$$M_{nt}(u) = \sum_{ij} \sigma_{ij} n_i t_j \quad (\text{twisting moment})$$

$$T(u) = \sigma_{ij,j} n_i + M_{nt,t} \quad (\text{normal shear force}).$$

Let $K \subset \mathbb{R}^2$ be a polygonal domain and set

$$a^K(u, v) := D \int_K \nu \Delta u \Delta v + (1 - \nu)(u_{,ij} v_{,ij}) dx.$$

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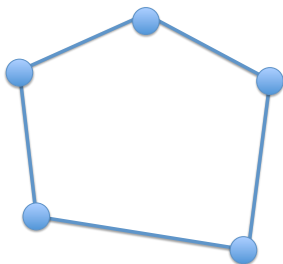
Integrations by parts yield the following useful equalities

$$\begin{aligned} a^K(u, v) &= D \left\{ \int_K \Delta^2 u v dx + \int_{\partial K} (\Delta u - (1 - \nu) u_{,tt}) v_{,n} ds \right. \\ &\quad \left. - \int_{\partial K} (\partial_n(\Delta u) v - (1 - \nu) u_{,nt} v_{,t}) ds \right\} \\ &= D \left\{ \int_K \Delta^2 u v dx + \int_{\partial K} M_{nn}(u) \partial_n v - \int_{\partial K} T(u) v ds \right. \\ &\quad \left. - \sum_{e \in \partial K} (M_{nt}(u), v n_{\partial e})_{\partial e} \right\}. \end{aligned}$$

where ∂e is the boundary of the edge e , $n_{\partial e}$ is the outwards normal to ∂e .

Towards nonconforming VEM discretization

- \mathcal{T}_h decomposition of Ω into polygons K



- h_K is the diameter of K , i.e., the maximum distance between any two vertices of K .
- On \mathcal{T}_h we make the following assumptions (for analysis only):

there exists a fixed number $\rho_0 > 0$ independent of \mathcal{T}_h , such that for every element K (with diameter h_K) it holds

- 1 K is star-shaped with respect to all the points of a ball of radius $\rho_0 h_K$
- 2 every edge $e \in \mathcal{E}_h$ has length $|e| \geq \rho_0 h_K$

Low-order nonconforming VEM

Local nonconforming VEM space:

$$V_{h,2}^K = \{v_h \in H^2(K) : \Delta^2 v_h = 0, M_{nn}(v_h)|_e \in \mathbb{P}^0(e), T(v_h)|_e = 0 \quad \forall e \in \partial K\}$$

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The **DOFs** of $V_{h,2}^K$ are chosen as follows:

- (d1) $v_h(\nu_i)$ for any vertex ν_i of K ;
- (d2) $\int_e \partial_n v_h ds$ for any edge e of ∂K .

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Remarks:

- Dofs (d1)-(d2) are unisolvent for $V_{h,2}^K$
- $\mathbb{P}^2(K) \subset V_{h,2}^K$.
- In case of a triangular mesh, we recover Morley's DOFs.
- The solution of the biharmonic problem appearing in the definition of $V_{h,2}^K$ is uniquely defined up to a linear function (which is however filtered by fixing, e.g., the values of the function in three non-aligned vertexes).

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$$a^K(v_h, v_h) = \int_K \Delta^2 v_h v_h dx - \sum_{e \in \partial K} (M_{nt}(v_h), v_h n_{\partial e})_{\partial e} + \int_{\partial K} M_{nn}(v_h) \partial_n v_h ds$$

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Coercivity gives $v_h = 0$ (note that setting to zero the nodal values of v_h filters the linear polynomials)

The global non-conforming low-order VEM space is defined as

$$V_{h,2} := \left\{ v_h : v_h|_K \in V_{h,2}^K, v_h \text{ continuous at internal vertexes,} \right. \\ \left. \int_e [\partial_n v_h] ds = 0 \quad \forall e \in \mathcal{E}_h^i, \right. \\ \left. \int_e \partial_n v_h ds = 0 \quad \forall e \in \mathcal{E}_h^\Gamma, v_h(\nu_i) = 0 \quad \forall \nu_i \in \mathcal{V}_h^\Gamma \right\},$$

where $[\cdot]$ is the usual jump operator.

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Note that functions in V_h are not continuous.

High-order nonconforming VEM

For $\ell > 2$, the local VEM space is defined as follows:

$$V_{h,\ell}^K = \{v_h \in H^2(K) : \Delta^2 v_h \in \mathbb{P}^{\ell-4}(K), M_{nn}(v_h)|_e \in \mathbb{P}^{\ell-2}(e), \\ T(v_h) \in \mathbb{P}^{\ell-3}(e) \forall e \in \partial K\}$$

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The **local DOFs** are chosen as follows:

- (D1) $v_h(\nu_i)$ for any vertex ν_i of K ;
- (D2) $\frac{1}{|K|} \int_K v_h p ds$ for any $p \in \mathbb{P}^{\ell-4}(K)$;
- (D3) $\int_e \partial_n v_h q ds$ for any $q \in \mathbb{P}^{\ell-2}(e)$ and any edge e of ∂K ;
- (D4) $\frac{1}{|e|} \int_e v_h q ds$ for any $q \in \mathbb{P}^{\ell-3}(e)$ and any edge e of ∂K .

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Remarks:

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$$\begin{aligned}
 V_{h,\ell} := & \left\{ v_h : v_h|_K \in V_h^K, v_h \text{ continuous at internal vertexes,} \right. \\
 & v_h(\nu_i) = 0 \quad \forall \nu_i \in \mathcal{V}_h^\Gamma \\
 & \int_e [\partial_n v_h] q ds = 0 \quad \forall q \in \mathbb{P}^{\ell-2}(e), \int_e [v_h] p ds = 0 \quad \forall p \in \mathbb{P}^{\ell-3}(e) \quad \forall e \in \mathcal{E}_h^i; \\
 & \left. \int_e \partial_n v_h q ds = 0 \quad \forall q \in \mathbb{P}^{\ell-2}(e), \int_e v_h p ds = 0 \quad \forall p \in \mathbb{P}^{\ell-3}(e) \quad \forall e \in \mathcal{E}_h^\Gamma \right\}.
 \end{aligned}$$

Local discrete bilinear form

We build a discrete local bilinear form $a_h^K(\cdot, \cdot) : V_{h,\ell}^K \times V_{h,\ell}^K \rightarrow \mathbb{R}$ such that

- (**ℓ -consistency**) $\forall p \in \mathbb{P}^\ell(K) \quad \forall v_h \in V_{h,\ell}^K$

$$a_h^K(p, v_h) = a^K(p, v_h)$$

- (**stability**) There exist positive constants α_*, α^* independent of h and K such that

$$\alpha_* a_h^K(v_h, v_h) \leq a_h^K(v_h, v_h) \leq \alpha^* a^K(v_h, v_h) \quad \forall v_h \in V_{h,\ell}^K.$$

- (**computability**) $\forall u_h, v_h \in V_{h,\ell}^K$ the local bilinear form

$$a_h^K(u_h, v_h)$$

is computable using only the DOFs of $V_{h,\ell}^K$.

Discrete VEM problem

Assemble in the usual way the global discrete bilinear form $a_h(\cdot, \cdot)$, i.e.

$$a_h(\cdot, \cdot) = \sum_{K \in \mathcal{T}_h} a_h^K(\cdot, \cdot).$$

The VEM discrete problem reads as follows: find $u_h \in V_{h,\ell}$ such that

$$a_h(u_h, v_h) = \langle f_h, v_h \rangle$$

for any $v_h \in V_{h,\ell}$, where $\langle f_h, v_h \rangle \simeq F(v_h)$.

Theorem (Antonietti, Manzini, V. 2016)

Under regularity assumptions on the mesh, there exists a unique solution $u_h \in V_{h,\ell}$. Moreover, it holds

$$|u - u_h|_{2,h} \lesssim h^{\ell-1} + \sup_{v_h \in V_{h,\ell}} \frac{\langle f - f_h, v_h \rangle}{|v_h|_{2,h}} + \sup_{v_h \in V_{h,\ell}} \frac{\mathcal{N}(u, v_h)}{|v_h|_{2,h}}$$

where

$$\mathcal{N}(u, v_h) := D \sum_{K \in \mathcal{T}_h} \left\{ \int_{\partial K} (\Delta u - (1 - \nu)u_{,tt}) v_{h,n} ds - \int_{\partial K} (\partial_n(\Delta u) v_h - (1 - \nu)u_{,nt} v_{h,t}) ds \right\}.$$

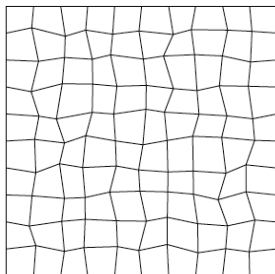
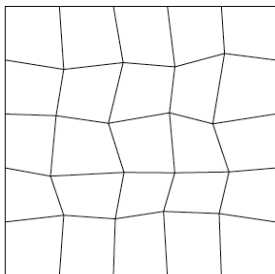
Corollary

Under regularity assumption on the mesh, the unique solution $u_h \in V_{h,\ell}$ satisfies the following error estimate, for $\ell \geq 2$

$$|u - u_h|_{2,h} \lesssim h^{\ell-1}.$$

Numerical results

- Exact solution: $u(x, y) = x^2(1 - x)^2y^2(1 - y)^2$
- Examples of employed polygonal meshes



Case $\ell = 2$

8 DOFs for each polygon

n	#dofs	Resid	$\mathcal{E}_{abs}(u_h)$	$\mathcal{E}_{rel}(u_h)$	Rate
0	$9.600 \cdot 10^1$	$3.823 \cdot 10^{-18}$	$6.671 \cdot 10^{-2}$	$9.364 \cdot 10^{-1}$	--
1	$3.410 \cdot 10^2$	$4.663 \cdot 10^{-18}$	$3.897 \cdot 10^{-2}$	$4.969 \cdot 10^{-1}$	0.999
2	$1.281 \cdot 10^3$	$7.007 \cdot 10^{-18}$	$2.053 \cdot 10^{-2}$	$2.559 \cdot 10^{-1}$	1.002
3	$4.961 \cdot 10^3$	$1.636 \cdot 10^{-17}$	$1.038 \cdot 10^{-2}$	$1.287 \cdot 10^{-1}$	1.015
4	$1.952 \cdot 10^4$	$3.694 \cdot 10^{-17}$	$5.224 \cdot 10^{-3}$	$6.467 \cdot 10^{-2}$	1.004
5	$7.744 \cdot 10^4$	$7.052 \cdot 10^{-17}$	$2.616 \cdot 10^{-3}$	$3.237 \cdot 10^{-2}$	1.004
6	$3.085 \cdot 10^5$	$1.771 \cdot 10^{-16}$	$1.309 \cdot 10^{-3}$	$1.620 \cdot 10^{-2}$	1

Case $\ell = 3$

16 DOFs for each polygon

n	#dofs	Resid	$\mathcal{E}_{abs}(u_h)$	$\mathcal{E}_{rel}(u_h)$	Rate
0	$2.160 \cdot 10^2$	$5.022 \cdot 10^{-18}$	$2.253 \cdot 10^{-2}$	$2.807 \cdot 10^{-1}$	--
1	$7.810 \cdot 10^2$	$9.850 \cdot 10^{-18}$	$7.419 \cdot 10^{-3}$	$9.185 \cdot 10^{-2}$	1.738
2	$2.961 \cdot 10^3$	$2.350 \cdot 10^{-17}$	$2.008 \cdot 10^{-3}$	$2.485 \cdot 10^{-2}$	1.961
3	$1.152 \cdot 10^4$	$4.610 \cdot 10^{-17}$	$5.214 \cdot 10^{-4}$	$6.452 \cdot 10^{-3}$	1.985
4	$4.544 \cdot 10^4$	$1.028 \cdot 10^{-16}$	$1.322 \cdot 10^{-4}$	$1.636 \cdot 10^{-3}$	2
5	$1.805 \cdot 10^5$	$2.122 \cdot 10^{-16}$	$3.330 \cdot 10^{-5}$	$4.121 \cdot 10^{-4}$	1.998
6	$7.194 \cdot 10^5$	$5.766 \cdot 10^{-16}$	$8.341 \cdot 10^{-6}$	$1.032 \cdot 10^{-4}$	2.002

Case $\ell = 4$

25 DOFs for each polygon

n	#dofs	Resid	$\mathcal{E}_{abs}(u_h)$	$\mathcal{E}_{rel}(u_h)$	Rate
0	$3.610 \cdot 10^2$	$5.301 \cdot 10^{-18}$	$3.589 \cdot 10^{-3}$	$4.441 \cdot 10^{-2}$	—
1	$1.321 \cdot 10^3$	$5.054 \cdot 10^{-17}$	$6.776 \cdot 10^{-4}$	$8.385 \cdot 10^{-3}$	2.570
2	$5.041 \cdot 10^3$	$2.445 \cdot 10^{-17}$	$1.125 \cdot 10^{-4}$	$1.392 \cdot 10^{-3}$	2.681
3	$1.116 \cdot 10^4$	$5.485 \cdot 10^{-17}$	$3.693 \cdot 10^{-5}$	$4.569 \cdot 10^{-4}$	2.802
4	$1.968 \cdot 10^4$	$5.240 \cdot 10^{-17}$	$1.638 \cdot 10^{-5}$	$2.027 \cdot 10^{-4}$	2.865
5	$3.060 \cdot 10^4$	$8.506 \cdot 10^{-17}$	$8.633 \cdot 10^{-6}$	$1.068 \cdot 10^{-4}$	2.902
6	$4.392 \cdot 10^4$	$7.660 \cdot 10^{-17}$	$5.089 \cdot 10^{-6}$	$6.297 \cdot 10^{-5}$	2.925
7	$5.964 \cdot 10^4$	$1.076 \cdot 10^{-16}$	$3.245 \cdot 10^{-6}$	$4.016 \cdot 10^{-5}$	2.940

Case $\ell = 5$

35 DOFs for each polygon

n	#dofs	Resid	$\mathcal{E}_{abs}(u_h)$	$\mathcal{E}_{rel}(u_h)$	Rate
0	$5.310 \cdot 10^2$	$2.115 \cdot 10^{-18}$	$1.605 \cdot 10^{-3}$	$1.985 \cdot 10^{-2}$	---
1	$1.961 \cdot 10^3$	$5.428 \cdot 10^{-17}$	$1.372 \cdot 10^{-4}$	$1.698 \cdot 10^{-3}$	3.764
2	$7.521 \cdot 10^3$	$1.077 \cdot 10^{-17}$	$8.904 \cdot 10^{-6}$	$1.102 \cdot 10^{-4}$	4.069
3	$1.668 \cdot 10^4$	$8.361 \cdot 10^{-17}$	$1.765 \cdot 10^{-6}$	$2.184 \cdot 10^{-5}$	4.063
4	$2.944 \cdot 10^4$	$2.374 \cdot 10^{-17}$	$5.595 \cdot 10^{-7}$	$6.924 \cdot 10^{-6}$	4.044
5	$4.580 \cdot 10^4$	$1.146 \cdot 10^{-16}$	$2.296 \cdot 10^{-7}$	$2.842 \cdot 10^{-6}$	4.030
6	$6.576 \cdot 10^4$	$3.491 \cdot 10^{-17}$	$1.109 \cdot 10^{-7}$	$1.373 \cdot 10^{-6}$	4.023
7	$8.932 \cdot 10^4$	$1.648 \cdot 10^{-16}$	$6.142 \cdot 10^{-8}$	$7.600 \cdot 10^{-7}$	3.860

Conclusions

We presented:

- Nonconforming VEM based discretization for a biharmonic problem on polygonal meshes
- Energy norm a priori error estimates
- Numerical results

Future work:

- Build an efficient preconditioner

References

- P.F. Antonietti, G. Manzini, M. Verani, The fully nonconforming Virtual Element Method for biharmonic problems, arXiv:1611.08736.
- J. Zhao, S. Chen, B. Zhang, The nonconforming virtual element method for plate bending problems. Math. Models Methods Appl. Sci. (2016).
- C. Chinosi, L.D. Marini, Virtual Element Methods for fourth order problems: L^2 estimates, Comput. Math. Appl., (2016).
- F. Brezzi, L.D. Marini, Virtual Element Method for plate bending problems, Comput. Methods Appl. Mech. Engrg., (2013).